

Intersubband Transitions: Their Physics and Their future.

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Outline

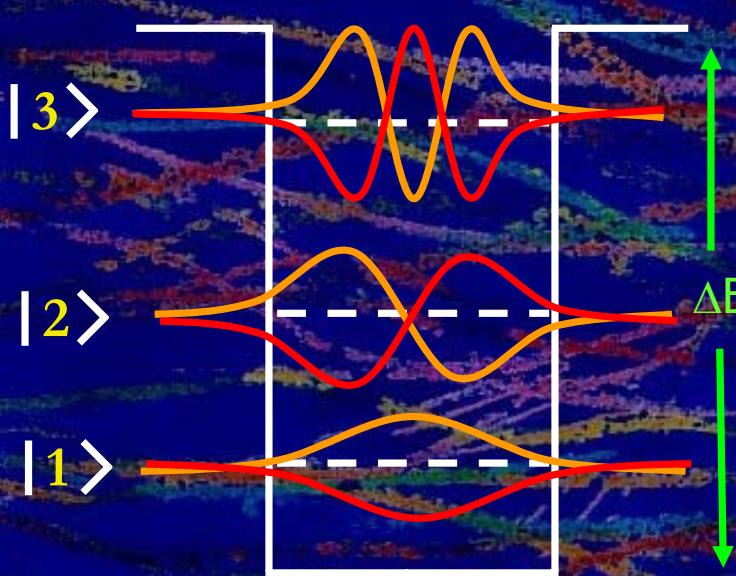
- *What are ISBT's?
- *What's good (and bad) about them?

BIG JUMP!

- *New wrinkles to try for improved devices.

- *New bits of science to try, just for the sake of it.

Confinement Energies



$$E_{12} \approx 3\hbar^2 / m^* d^2$$

$\Delta E_c * d \sim \text{few nm}$

$* \Delta E_c \sim \text{few 100 meV.}$

$* \text{Much wider range of material possibilities than interband-based devices.}$



100 years of living science



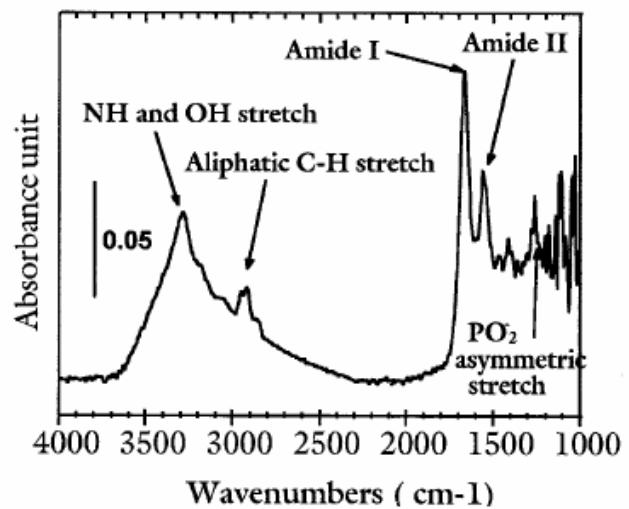
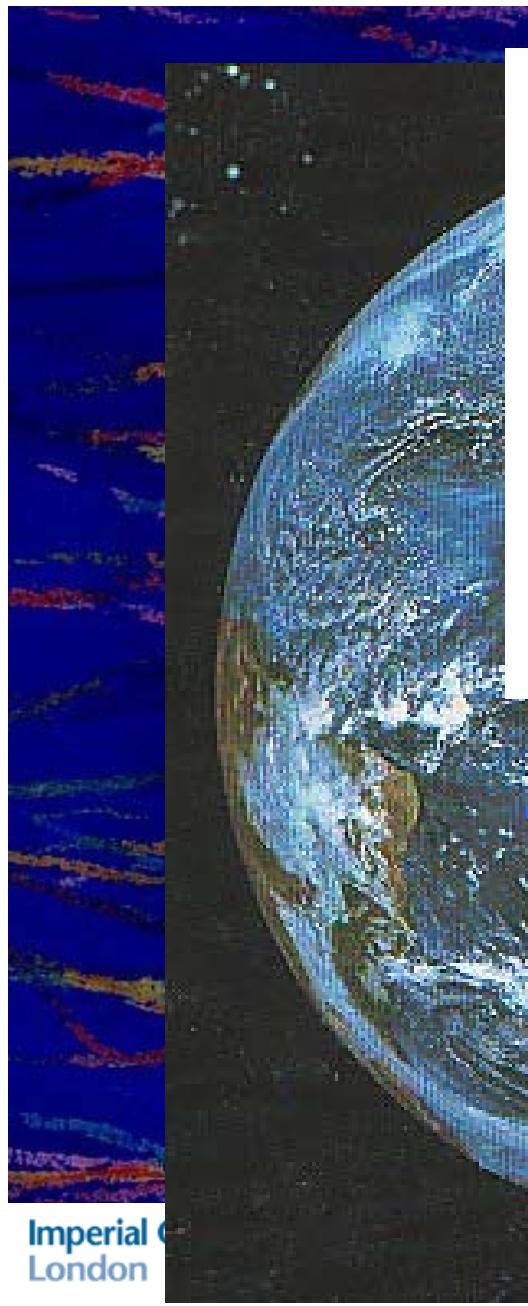
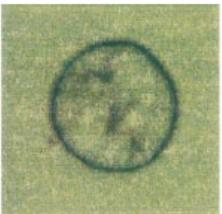


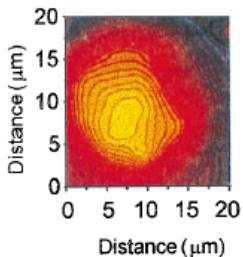
FIG. 1. Infrared spectrum of a mouse UN2 hybridoma B living cell, recorded with an aperture of $3 \times 3 \mu\text{m}^2$. The instrumental resolution was set at 4 cm^{-1} , and the spectrum displayed is the result of 128 co-added scans (the total recording time is 55 sec).

A.Optical Image

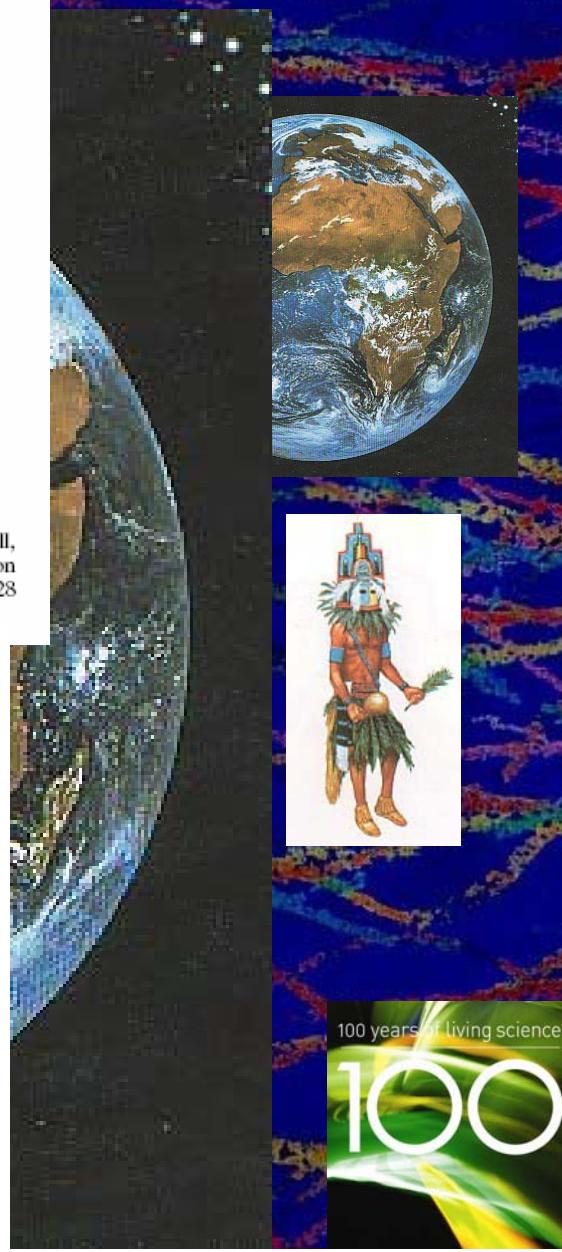
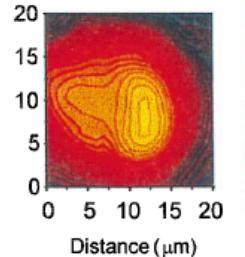
N Jamin et al PNAS
USA 95(1), 4837-4840
(1998)



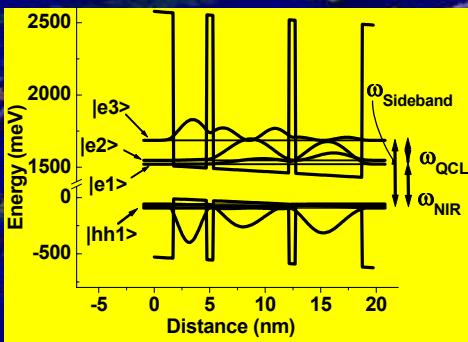
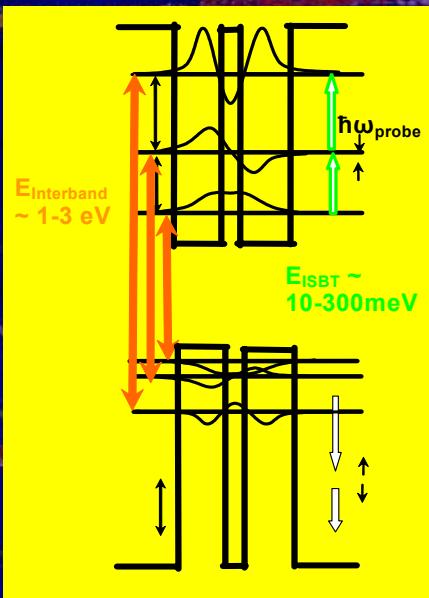
B.Proteins



C.Lipids



Transition strengths



* Δn -odd selection rule (in principle).

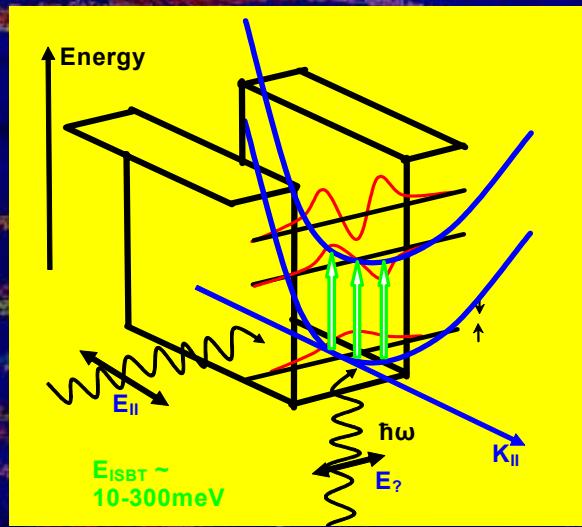
* $z_{12} \sim 2-3 \text{ nm}$. i.e. big!



* Highly anisotropic absorption (dichroism) and dispersion (birefringence).

* Optical response has unusual quasi-cylindrical point group symmetry.

Tunable ISBT oscillator strengths.



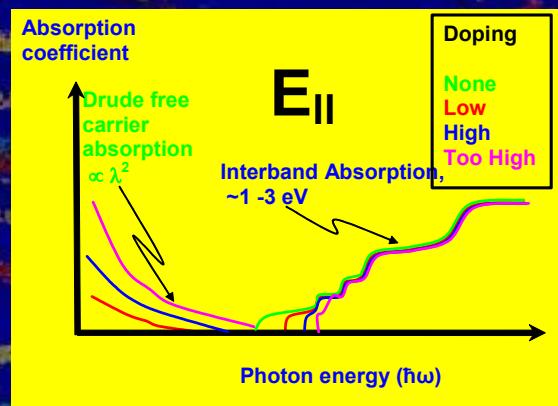
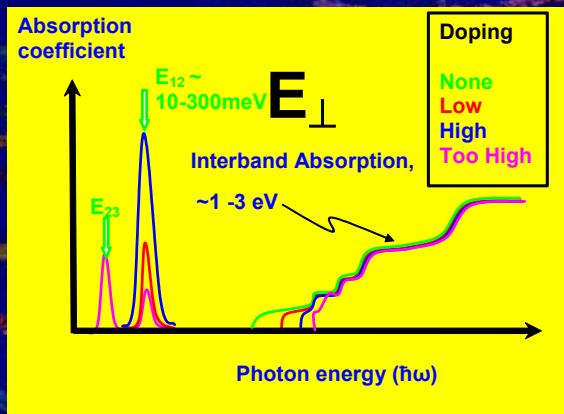
*ISBT transition energy (almost) independent of k_{\parallel}



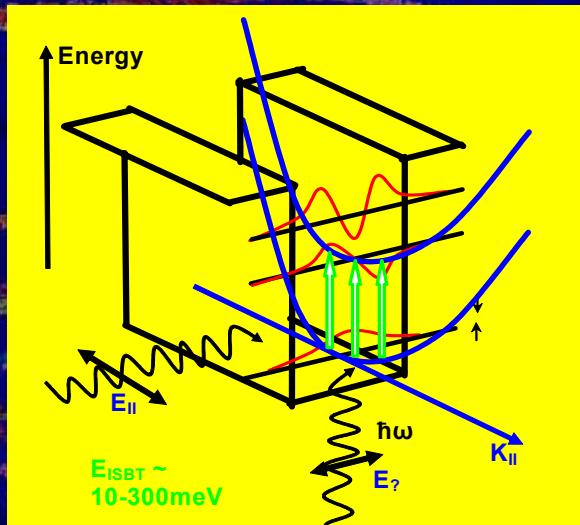
* ISBT absorption is peaked



* Absorption \propto electron density



Dipole selection rules.

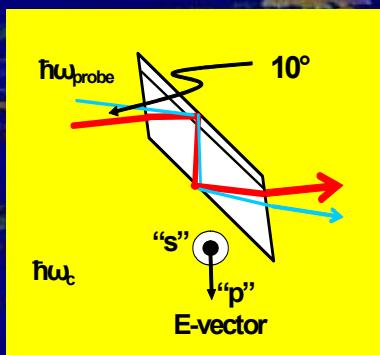


*ISBT oscillator is z-polarised



* Normally no absorption at normal incidence.

* No vertical laser cavities

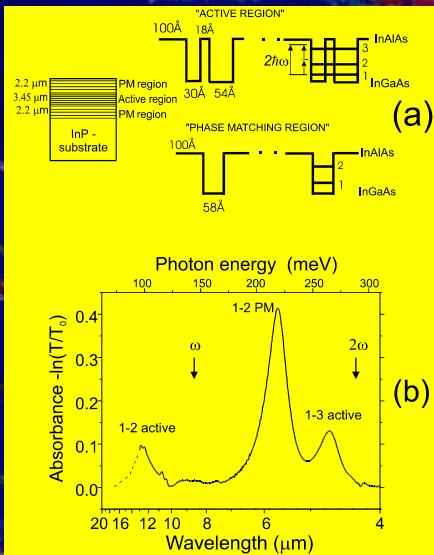


* Can be dodged with e.g VB and non- Γ CB material. In principle.



* Makes for easy absorption spectroscopy though!

Designable dispersion.

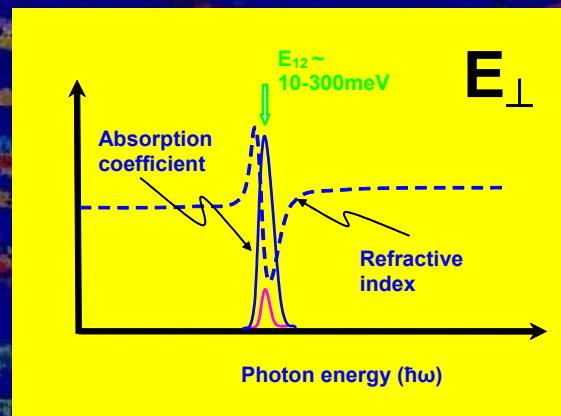


* ISBT peak is surrounded by transparent regions.

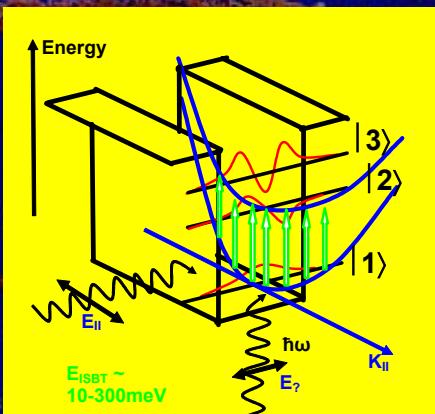
* Can add oscillators to waveguide to design-in desired dispersion characteristics.

* Useful for coherent optical experiments.

* Electro-optical control possible through QCSE.



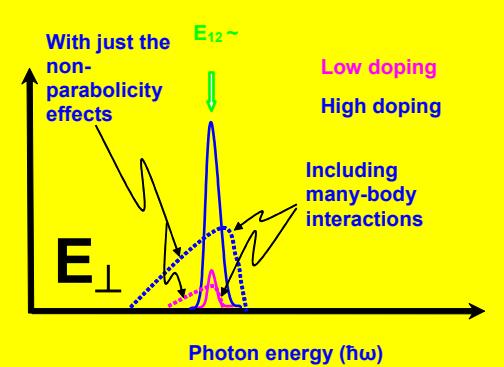
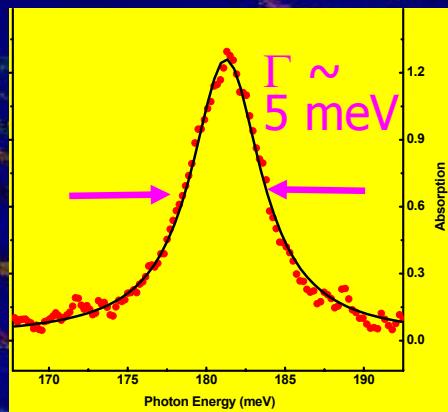
Broadening mechanisms (1).



* m^* increases with energy (non-parabolicity), broadens transition.

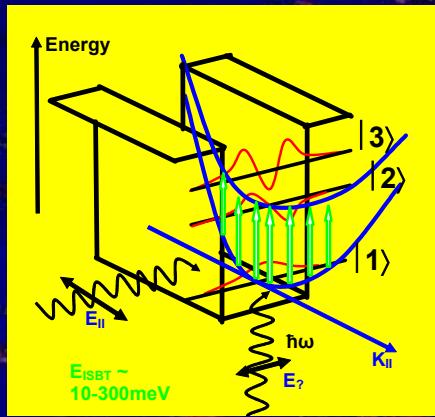
* Exchange correlation effects counteract this broadening effect.

* Often a good fit to Lorentzian lineshape.



Broadening mechanisms (2).

* 300k Linewidth ~ 10 meV , $<< kT$!

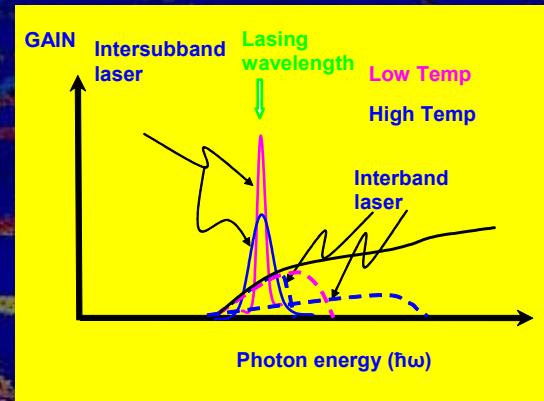
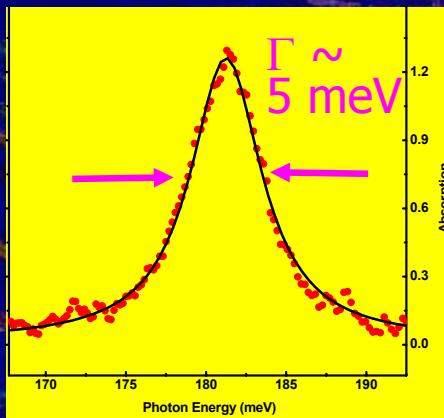


* Interface roughness and phonon scattering $W \sim 1/q^2$

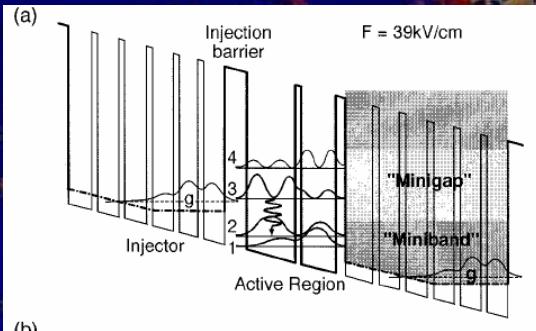
* Low-q acoustic Phonons do the damage it seems.

* Scattering respects dipole selection rules.

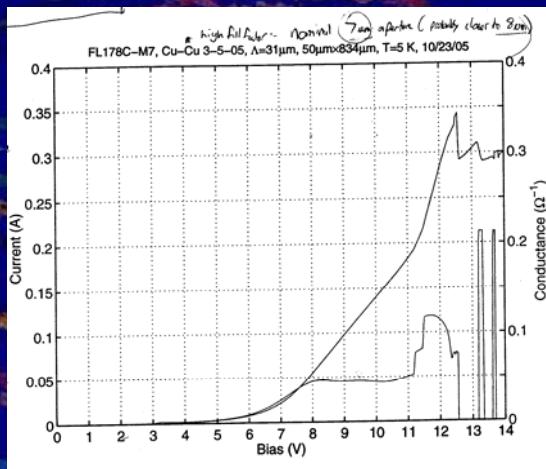
* Temperature insensitivity
VERY good for laser J_0 's



Cascaded Electron Transport Idea .



C Sirtori et al. APL 69 (19), 2810 (1996).



* Majority carriers, insensitive to Temp and Materials quality.

* Very versatile cf pn junctions.

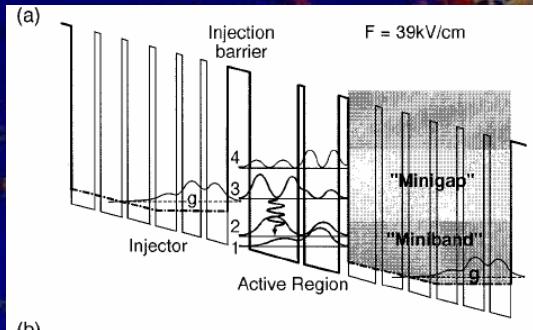
* Resonant tunnelling can help keep everything aligned.

* Suited to wide cavities needed for IR.

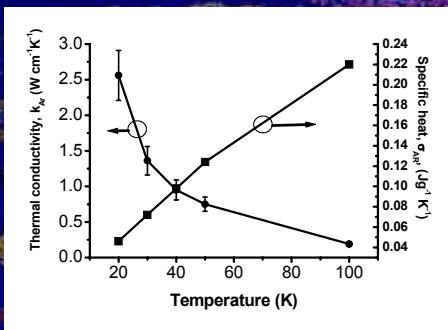
* Important "transformer" side effect.

Thermal Transport :

* It's rubbish. Frankly.



C Sirtori et al. APL 69 (19), 2810 (1996).



A J Borak et al. APL 82 (23), 4020 (2003).

* GaAs Phonon mfp's 1000 nm (50K) to $\sim 40\text{nm}$ (375K)

* Heterostructures scatter phonons on nm length scales

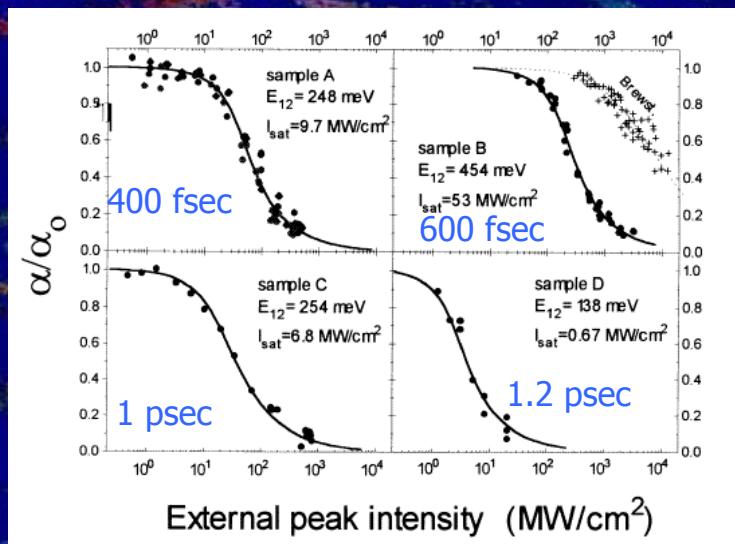
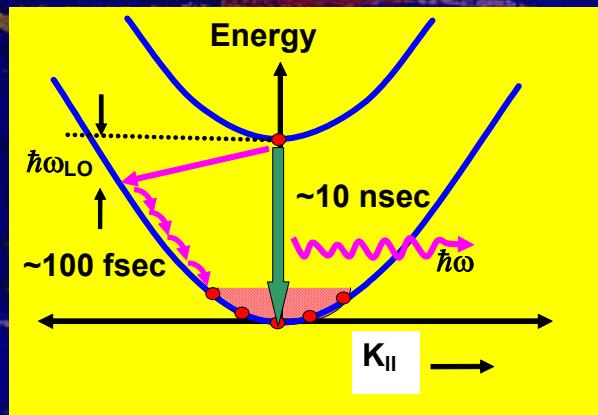
* QCL conductivities $\sim 10 \times$ worse than bulk

* Comparatively few systematic studies.

* CRITICAL for 300K CW devices.



Intersubband dynamics



K L Vodopyanov et al. Semicond. Sci. Tech !2 , 08
(1994)

* GaAs LO phonon scattering in $\sim 100 \text{ fsec}$.



* Faster than a very fast thing.



* Poor spontaneous radiative efficiency.



* High J_{th} in lasers



* High saturation intensities in switches.



* "Phonon bottleneck" failed to appear.



Half time summary .

* Intersubband transitions are.....

Tunable in energy and strength throughout IR and FIR

Fast, (due to intrinsic non-radiative recombination channel)

Quite sharply peaked ($\Delta\lambda/\lambda \sim 10\% @ 300K$).

Insensitive to temperature and materials.

Tunable in energy and strength throughout IR and FIR.

Fast, (due to intrinsic non-radiative recombination channel)

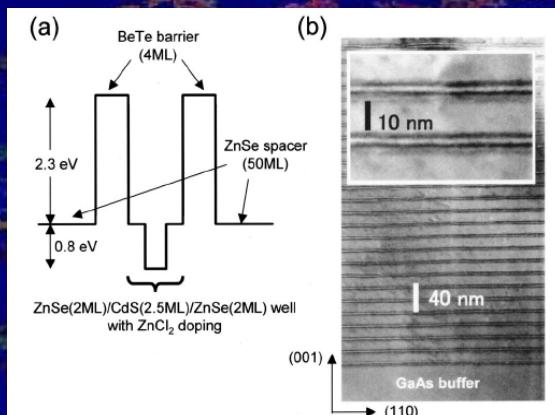
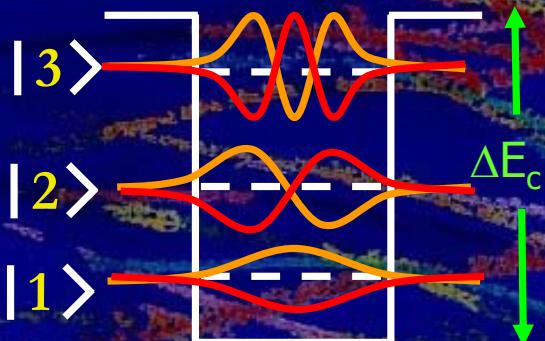
Quite sharply peaked ($\Delta\lambda/\lambda \sim 10\% @ 300K$).

Insensitive to temperature and materials.

What next then?

- * New Materials possibilities.
- * Plasmonics
- * Non-linear optics for new wavelengths.
- * Strongly coupled light-matter systems.
- * Coherent Quantum Optical effects.
- * Phonon Engineering.
- * Subbands for Science:-Artificial Atoms.

New Materials.

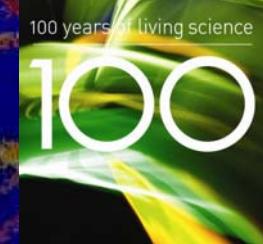
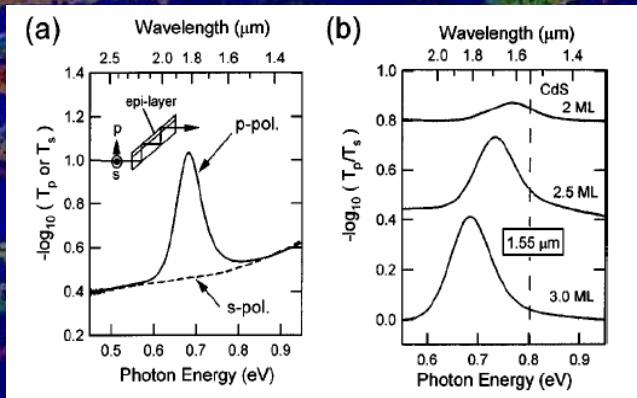


R Akimoto et al. APL 81(16), 2998 (2002).
[Tsukuba]

- * Shorter λ 's need larger ΔE_c materials
- * Larger ΔE_c goes with larger bandgap and m^* , and higher bond strengths.
- * This means higher temp growth and doping problems.

$$E_{12} \approx 3h^2 / m^* d^2$$

* Feeling Brave?



$$\omega_p^2 \approx ne^2 / m^* \epsilon_0 \epsilon_r$$

Plasmonics

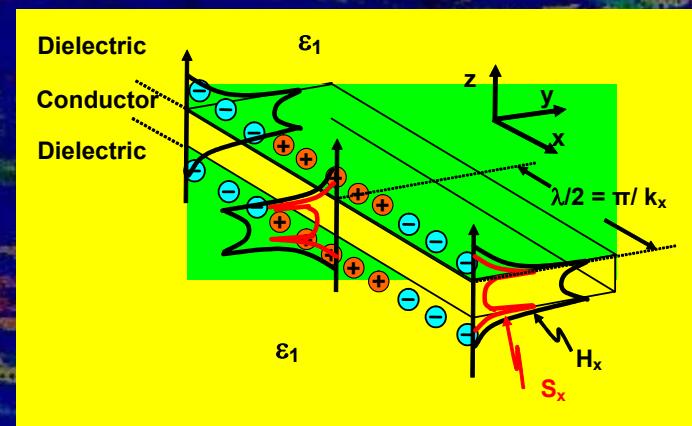
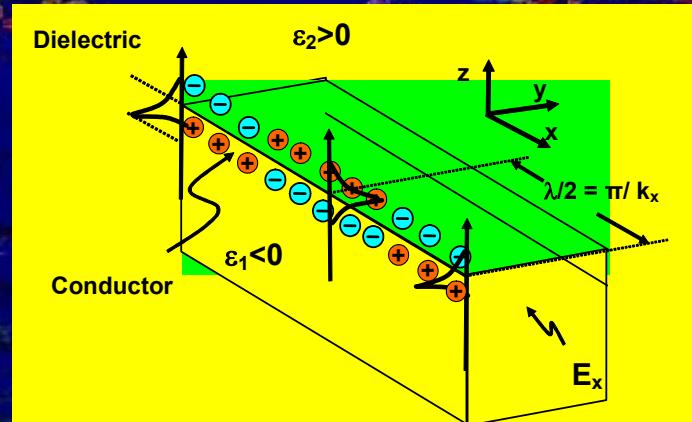
* $\epsilon < 0$ for $\omega < \omega_p$

* Field Concentration-sensors

* Shrinking Opto-Electronics

* Metallic Optics

* Opportunities for semiconductors!



Plasmonics

$$\omega_p^2 \approx n e^2 / m^* \epsilon_0 \epsilon_r$$

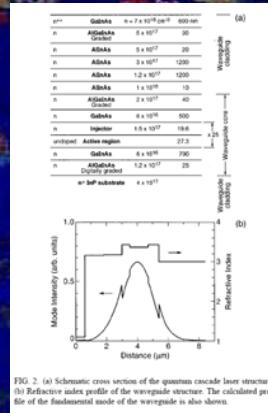
* Already being used for QCL cavities.

* Compresses mode volumes.

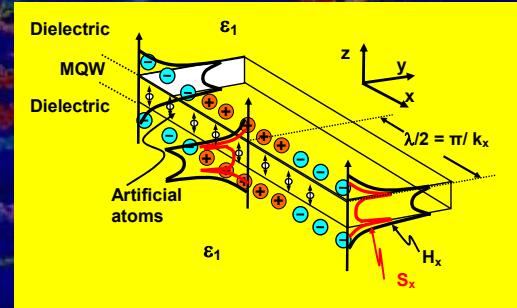
- Way of guiding FIR in sensors?

* Quantum anisotropy means new opportunities for semiconductors!

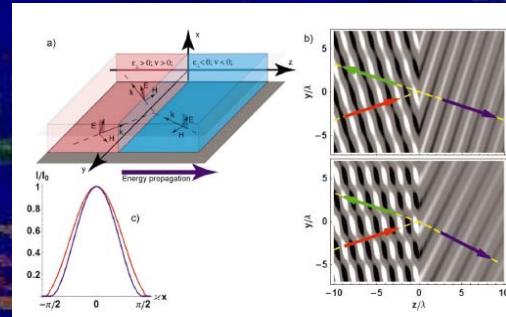
* Negative refraction exotica possible.



Sirtori et al. APL 66 (24), 3242 (1995).



Jonathan Plumridge and CCP. Phys. Rev. B 76, 075326 (Aug 2007).



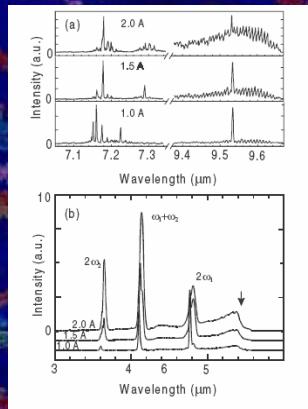
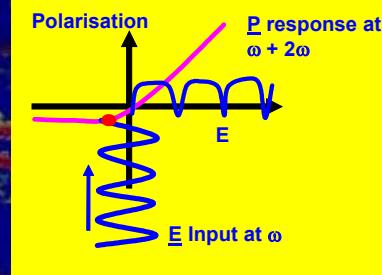
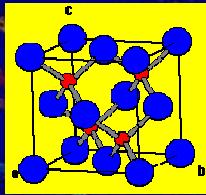
V.A. Podolskiy and E.E. Narimanov, PRB 71, 201101, (2005).

Non-linear optics for new wavelengths

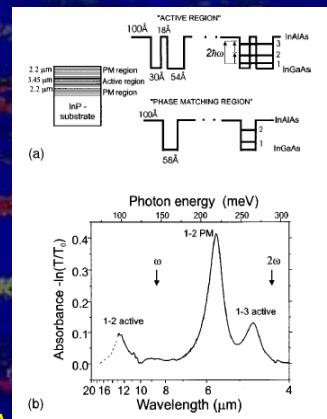
*QCL cavities have high radiation density and the intrinsic $\chi(2)$ of III-V's.

*Artificial resonances can be designed-in.

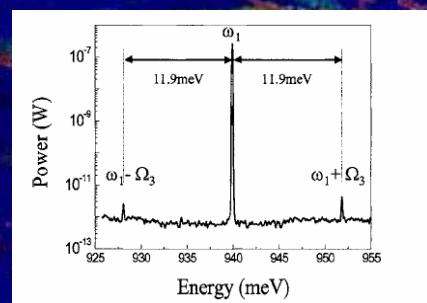
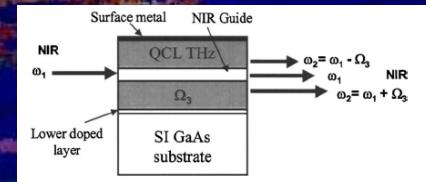
*Designable dispersion for phase matching.



N Owschinnikow et al. PRL, 90 (4), 043902-1 (2003).



ITQW Ambleside Sept 2007.



S S Dhillon et al. APL 87, 071101 (2005).

K L Vodopyanov et al. APL 72(21), 2654 (1998).

Strongly coupled cavities

*Einstein's rules assume matter and light only perturbatively coupled.

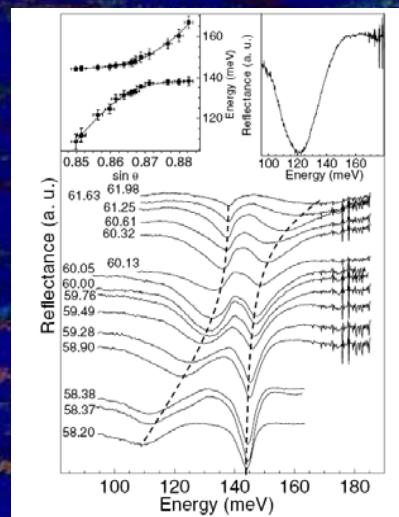
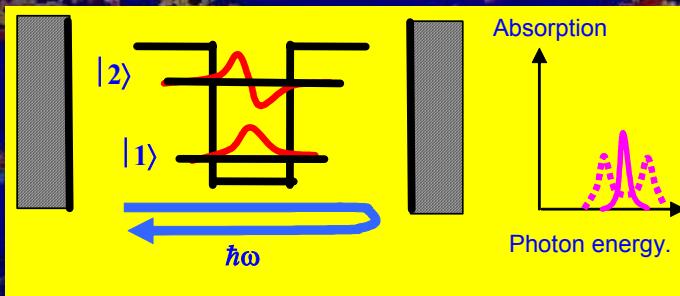
*Small mode volume, high Q, coupling energy dominates all other terms.

*Good for weakly absorbing detectors.

*Non-Einsteinian lasers?

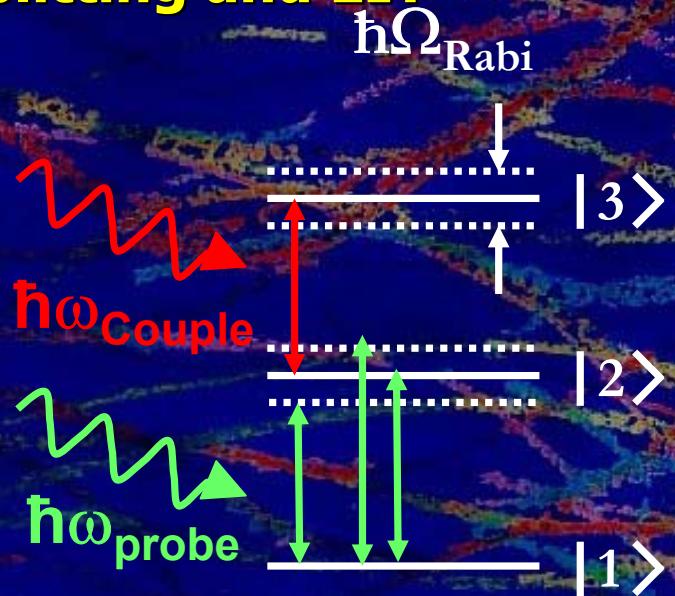
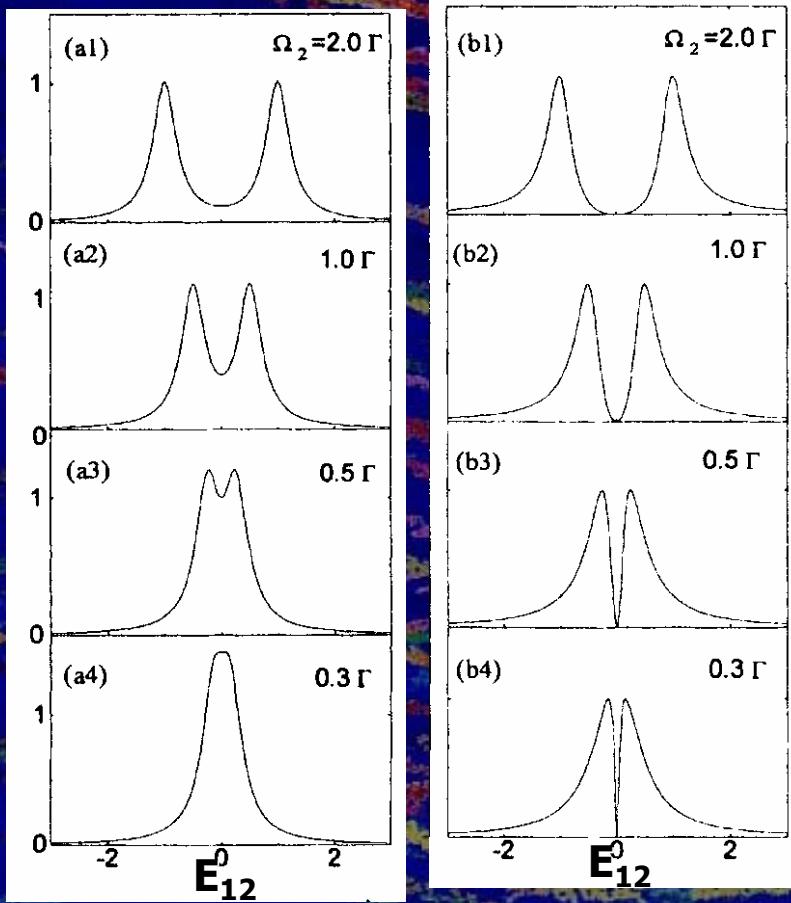
ET Jeynes and FW Cummings Proc, IEEE , 51, (1), 89-109 (1963).

*QCL's uniquely suitable for this.



D Dini et al. PRL, 90, 116401 (2003)

Quantum Optics: Rabi Splitting and EIT



Generalised Rabi Frequency
 $\hbar\Omega_{\text{Rabi}}^2 = (\mu \cdot E)^2 + (\hbar\Delta_{\text{Couple}})^2$

Just the dressing

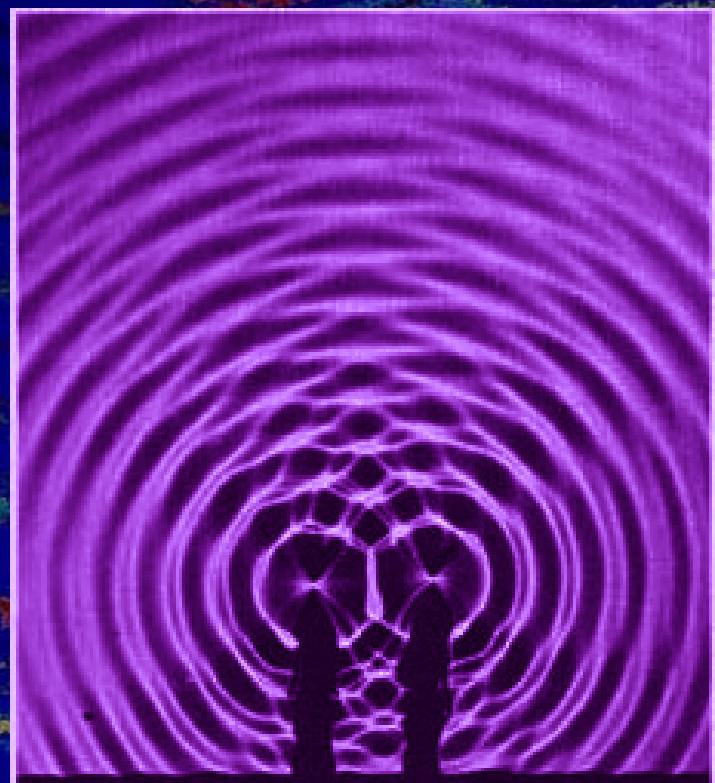
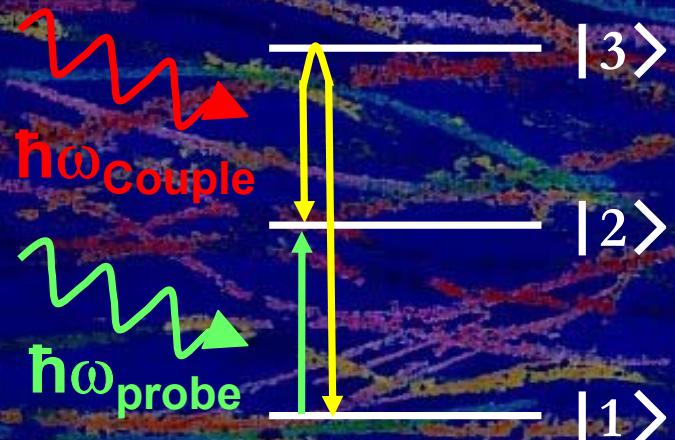
Imperial College
London

Dressing and
Quantum Interference

T1QW Ambleside Sept 2007.



Matter-wave interference



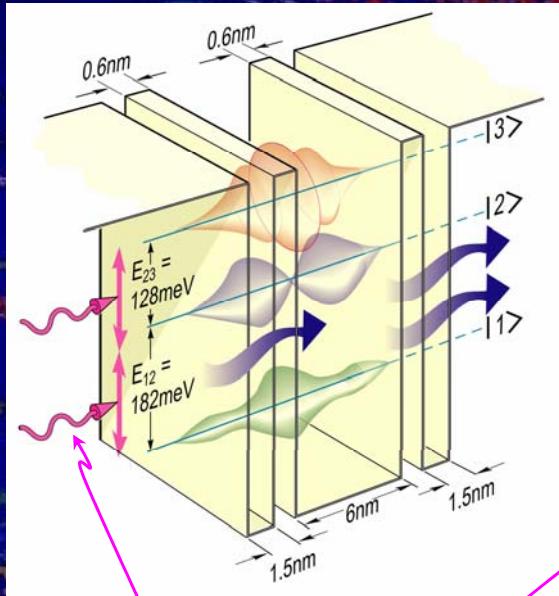
EIT was first seen in atomic vapours...



K J Boller, A Imamoglu and S E Harris, PRL
66,2593 (1991), using Rubidium vapour.

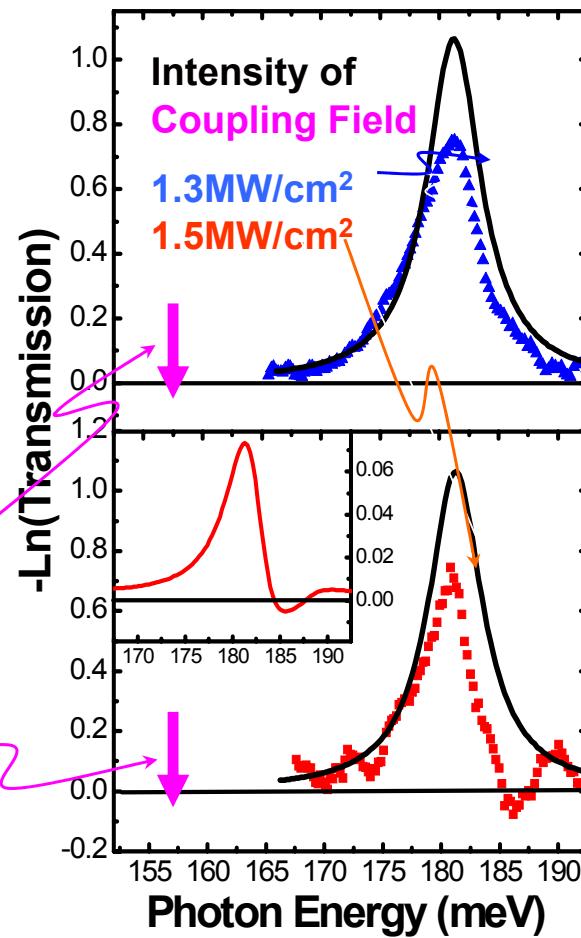
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GWI with subband transitions!

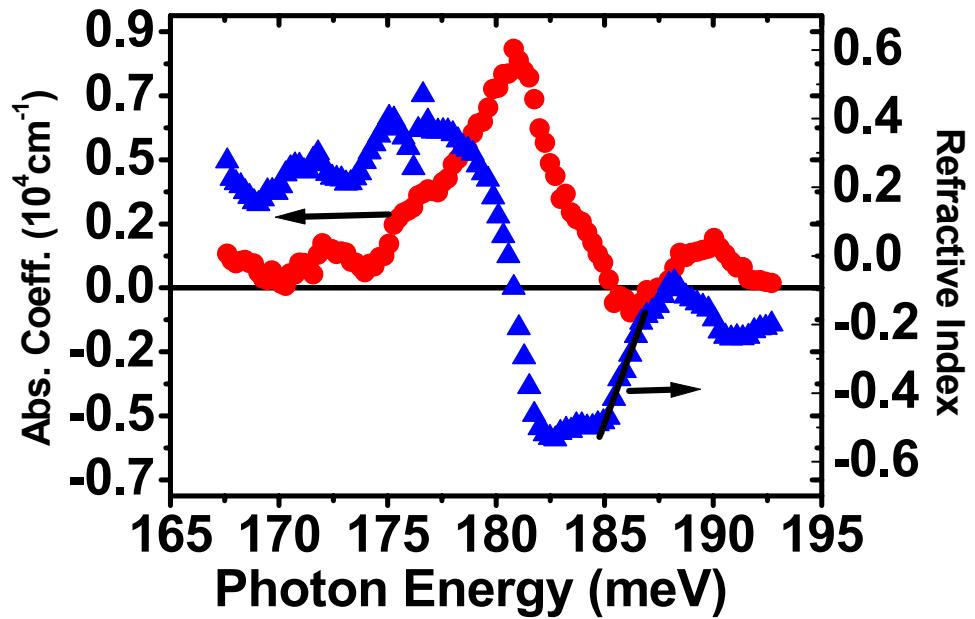


$$\hbar\omega_{\text{Couple}} = 155\text{meV}$$

M D Frogley et al, *Nature Materials*, Vol. 5, 175-178, March 2006.



Slow Light.



$$V_{\text{group}} = c/n_{\text{group}} = c/(n_0 + \omega d\mathbf{n}/d\omega)$$

$V_{\text{group}} = c/40$ in the region where the sample is transparent/amplifying.

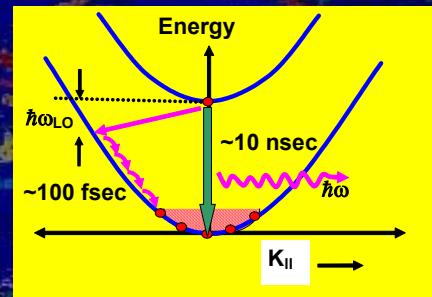
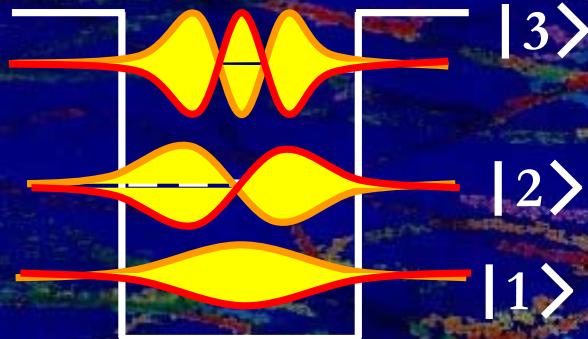
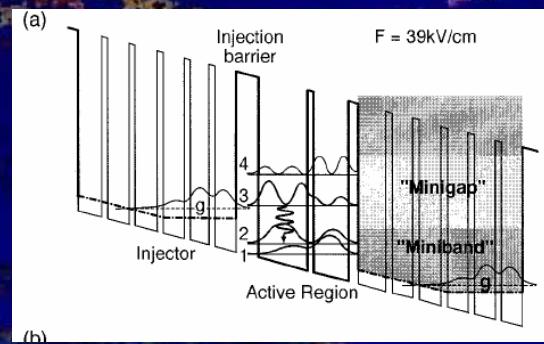
Phonon Engineering?

*We already structure electronics an optical aspects, why not acoustic as well?

*Dephasing Acoustic and depopulating phonons both have $\lambda \sim 15$ nm.

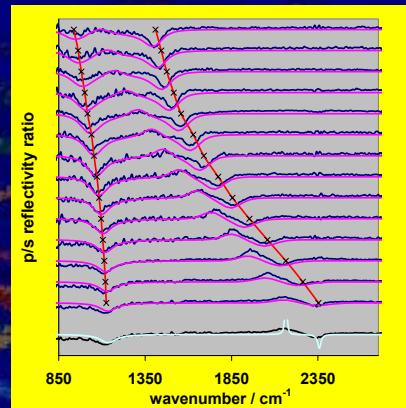
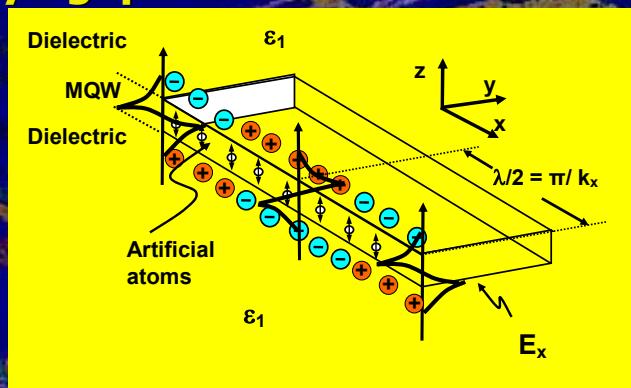
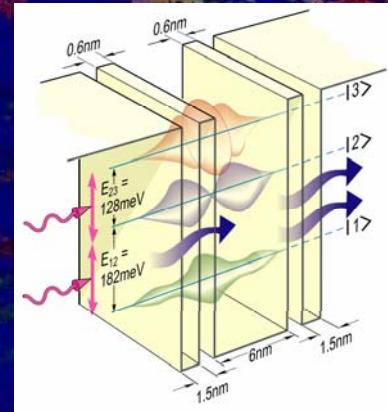
*Phononic SL's for improved thermal conductivity?

*Phononic anti-nodal cavities for decoupling electrons from crystal?



Artificial atoms for Science!

- *ISBT's behave like
designable atoms
- *New symmetries and energy
level juxtapositions possible.
- *Cavities are easy
- *They stay put!
- *Cavity QED?
- *Flying qubits?



Acknowledgements

Mark Frogley (ex-Imperial, now RAL)

Jon Plumridge (Imperial)

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Mary Matthews (Imperial)

Rob Steed (Imperial)

Ed Clarke (Imperial)

Ray Murray

H C Liu

Emmanuel Dupont

Ben Williams (ex-MIT, now UCLA)

Maurice Skolnick (Sheffield)

Luke Wilson (Sheffield)

Jerome Faist and
Mattias Beck (Neuchatel)

Elodie Roussard (Imperial)

Justin Rodger (Masters ICL)

Danny Segal, Almut Beige (ICL)

Spectroscopy
and modelling

MBE growth

Strongly coupled
Cavity samples

THz frequency mixing
QCL's

Artificial atom multilayers

DMT/Gaussian beam

AQW Design/DMT software

Discussion – Dressed states

EPSRC (UK) – Funding

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