Mode behavior, waveguide losses, and gain of

two-sectioned, coupled-cavity GaAs/(AI,Ga)As

THz and MIR QCLs

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Outline

1. Two-sectioned, coupled-cavity QCLs

Longitudinal modes observed for different:

- 2. *Lengths* of the two subcavities and analyzed within a transfer-matrix approach
- 3. *Current values* through the two subcavities
- 4. *Time delays* between both current pulses
- 5. Summary, outlook

T-sectioned, coupled-cavities (TSCC) lasers



Parameters to control TSCC lasers:

• total cavity length $L_{tot} = L_1 + L_3$

ratio between lengths L₃/L₁

• currents I_1 and I_3

MIR: $L_2 \leq$ laser wavelength THz: $L_2 <$ laser wavelength

Optical coupling between cavities *j* = 1 and 3 modifies longitudinal modes, indicated 1983 for interband lasers.

TSCC QCLs



We investigate longitudinal modes of MIR/THz TSCC QCLs,

 $L_2 \approx 0.4$ to 1 µm (Gap formed by cleaving the ridge.)

In the following we focus on results of THz devices

Length ratio of subcavities L_3/L_1 – mode suppression

THz QCLs: design according to Barbieri et al., APL 2004, single plasmon waveguides



 $I_{\rm tot} = I_1 + I_3 \approx 1.5 I_{\rm thr} \approx 0.9 I_{\rm max}$

Length ratio of subcavities L_3/L_1 – mode suppression



Model for mode suppression – transfer-matrix approach

- Photon emission does not depend on laser field: optical constants are independent of intensity
- Interference partial waves: transfer matrix
- No incoming light: *m* modes with complex eigenfrequencies $\hat{\nu}_m$ -

$$B_{0} \leftarrow \begin{bmatrix} A_{1}^{'} & A_{1} & A_{1}^{'} & A_{3}^{'} & j = 3 & A_{3}^{'} \\ B_{1}^{'} & j = 1 & B_{1}^{'} & B_{3}^{'} & j = 3 & B_{3}^{'} \end{bmatrix} \rightarrow A_{4}^{'}$$

$$egin{pmatrix} A_0 = 0 \ B_0 \end{pmatrix} = \mathrm{T} egin{pmatrix} A'_4 \ B'_4 = 0 \end{pmatrix}$$

$$\rightarrow \mathrm{T}_{11}(\hat{\nu}_m) = 0$$

Mode frequencies $u_m = {
m R}e\{\hat{
u}_m\}$, mode heights $-1/{
m I}m\{\hat{
u}_m\}$

Spectral function TSCC resonator $P(
u) = \mid T_{22} + r_{
m f}T_{21} - r_{
m b}T_{12} - r_{
m f}r_{
m b}T_{11}\mid^2$

Observed modes – calc. eigenfrequencies and spectral function



 $I_{\rm tot}$ =1.5 $I_{\rm thr}$ \approx 0.9 $I_{\rm max}$

Mode:

- suppression caused by multiplicative superposition of partial waves
- spacings ∆v described by T-matrix approach
- height modulation $\Delta S/S_0(\hat{\nu} < \Delta S/S_0(P))$ exp. $\Delta S/S_0 > \Delta S/S_0$ calc., stimulated emission not fully included in model
- ΔS related with Δν,
 Kramers-Kronig relation

Currents I_1 and I_3 – optical saturation and subband misalignment

both cavities exhibit lasing; $I_{max} \approx 0.9 \text{ A} \approx 2 I_{thr}$; $L_3/L_1 \approx 1$





 $I_{tot} < I_{max}$; S < 60; mode suppression

 $I_{tot} \approx I_{max}; S_{max} \approx 80$ • saturation laser output

- suppression disappears
- subband misalignment

 $I_{tot} > I_{max}$; S < 60; mode suppression

Optical saturation and subband misalignment, $L_1/L_3 \approx 1$





Low laser output, $S \approx 6 < S_{max} \approx 80$,

- no intensity saturation
- always mode suppression

Time delays δt between current pulses I_1 and I_3

 $L_1/L_3 \approx 1$; $I_1 = I_3 = 0.7 \text{ A} \approx 1.6 I_{\text{thr}} \approx 0.84 I_{\text{max}}$



Observed features longitudinal modes: described by T-matrix:

- suppression of $N = L_3/L_1$ modes interference effects \checkmark
- suppression disappears optical saturation –
- switching mode features dynamical effects?
- THz: laser field always in both subcavities
 ✓
 MIR: light can be focussed in one cavity

Stimulated emission

Interplay between:

photon emission, photon density, and electron populations has to be fully included in T-matrix approach