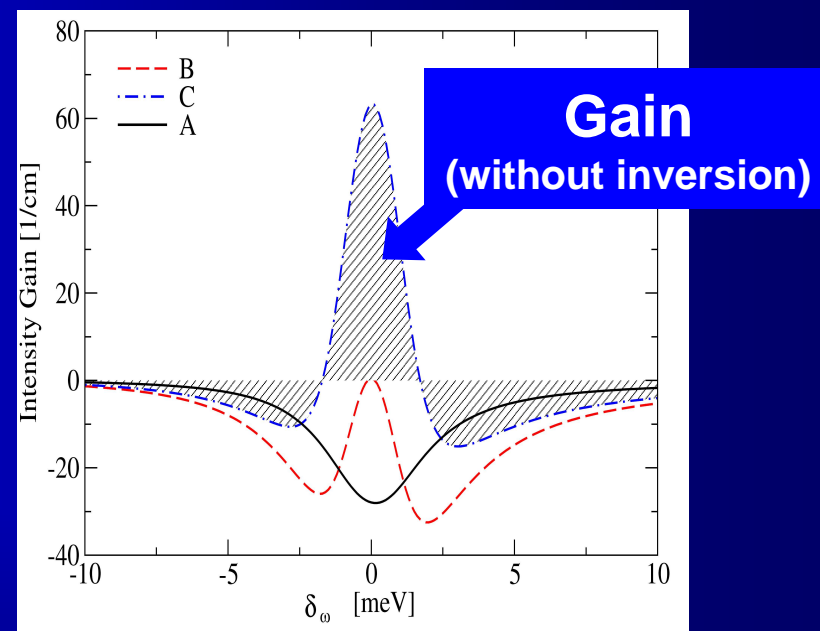


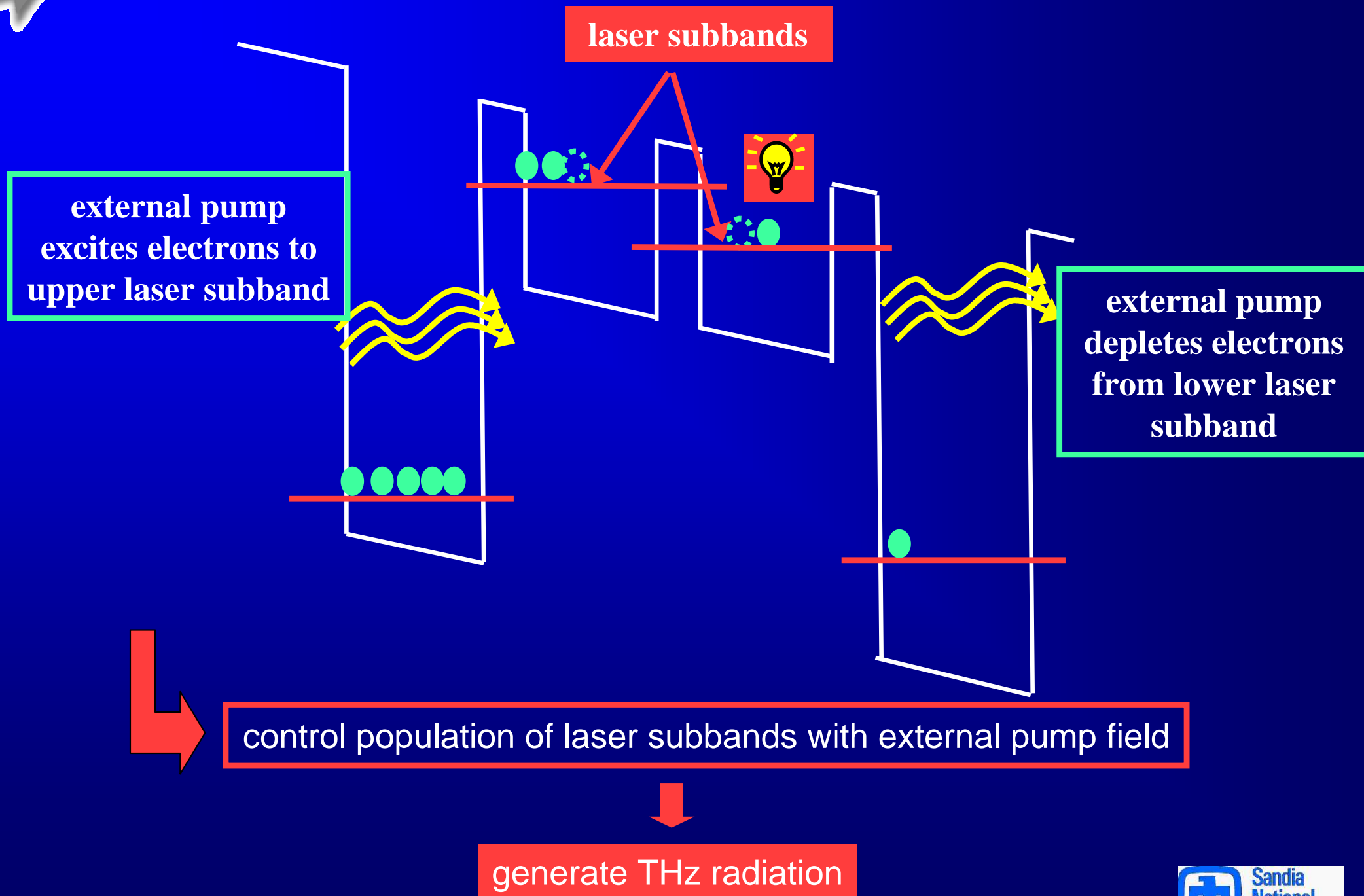
Who needs population inversion?

Automatically phase-matched
quantum coherence contributions
as a source for THz radiation



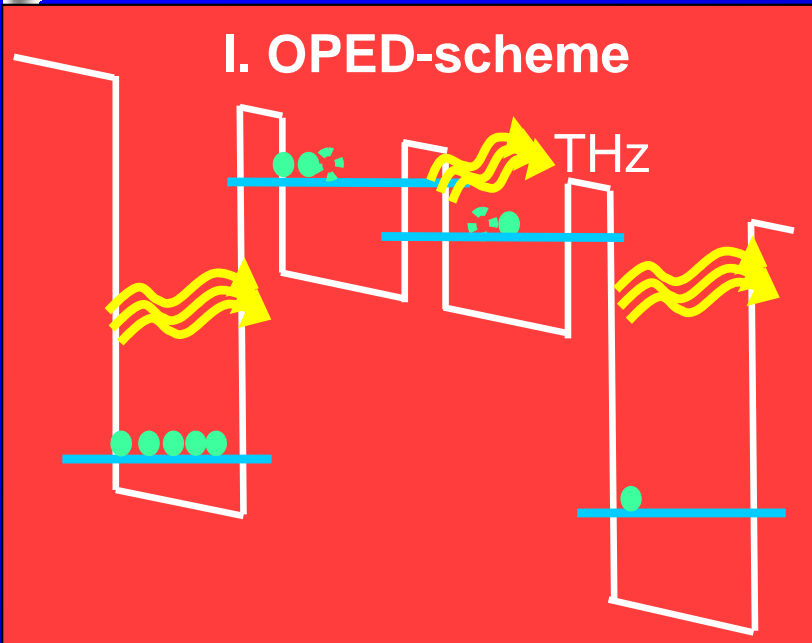
Inès Waldmueller, Weng W. Chow and Michael C. Wanke
Sandia National Laboratories

Opt. pumped, electrically driven QCL (OPED)

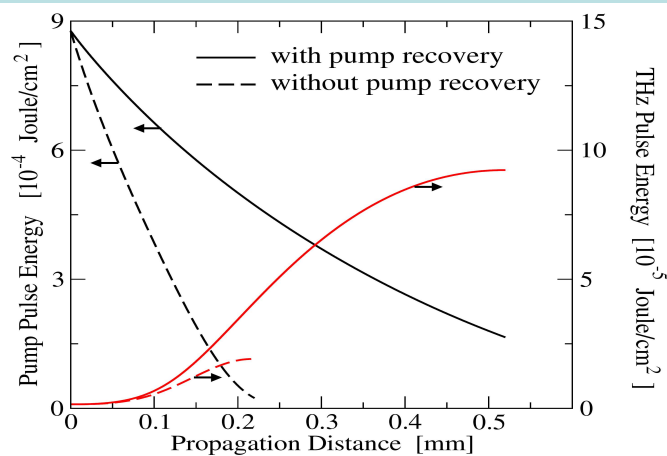


Overview

I. OPED-scheme

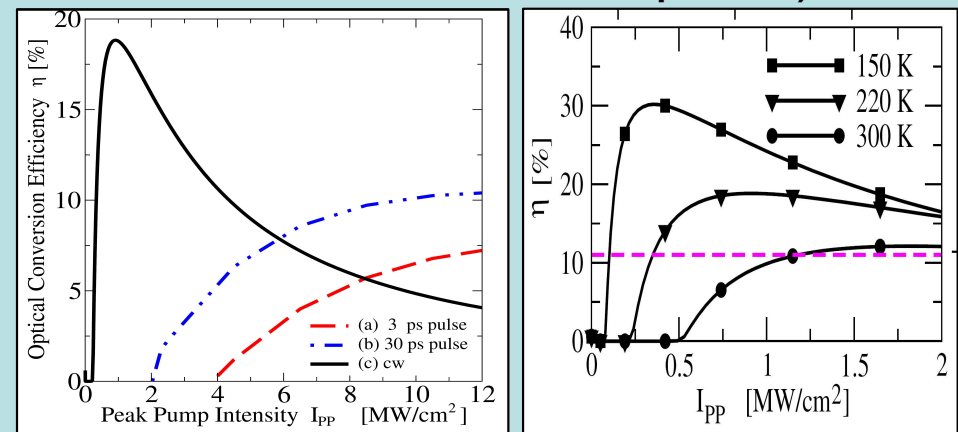


II. Benefits of Pump Recovery

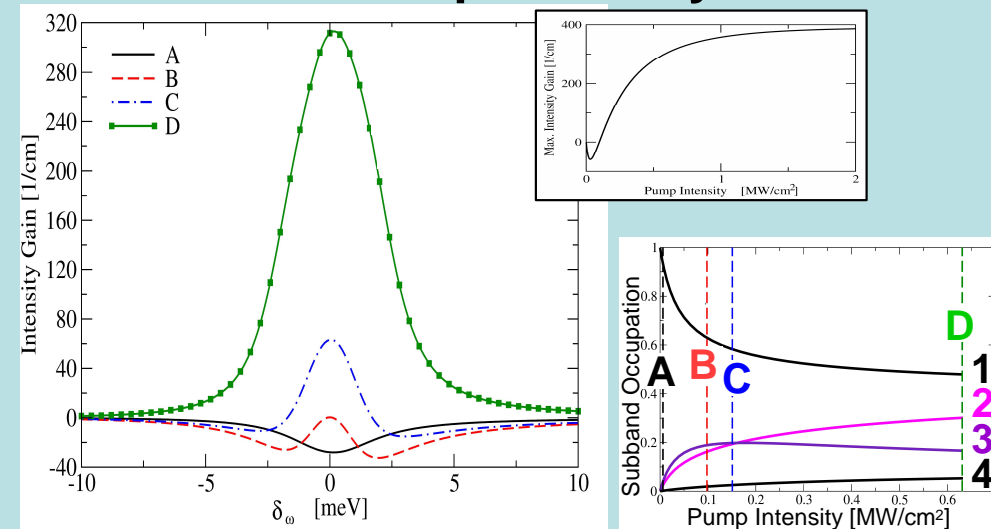


III. Capabilities of OPED

(Dependence of Optical Conversion Efficiency on Pulse Duration and Temperature)

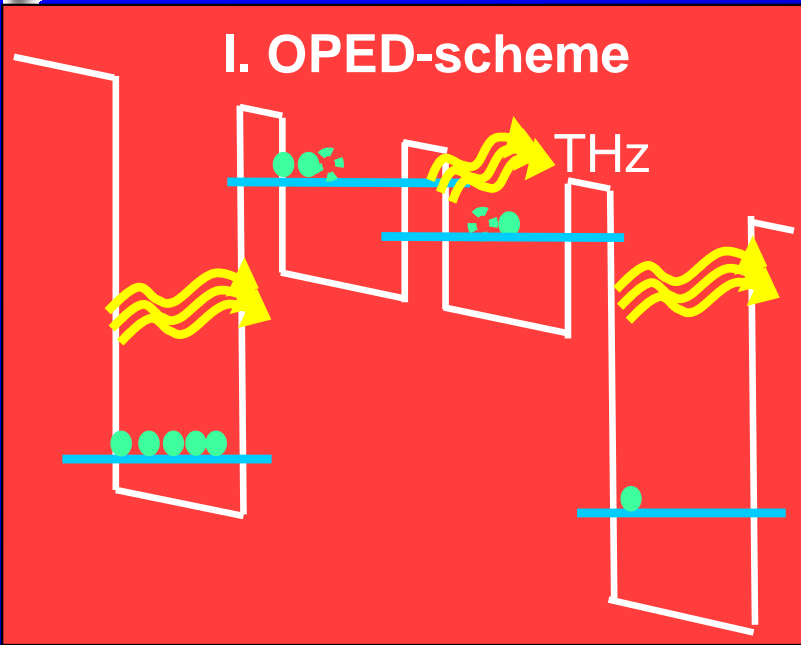


IV. Dependence of Small Signal Gain on Pump Intensity

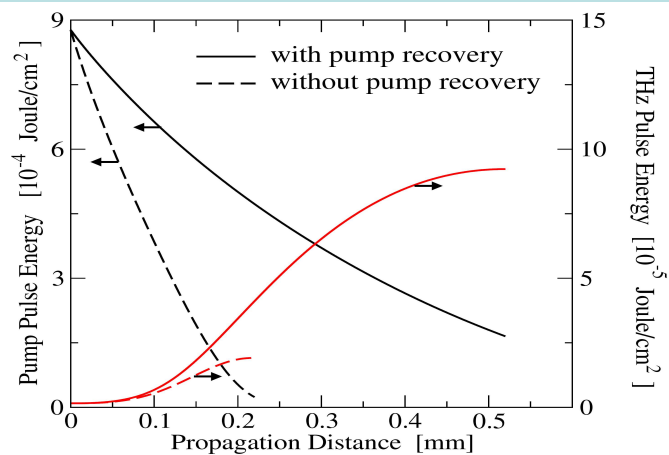


Overview

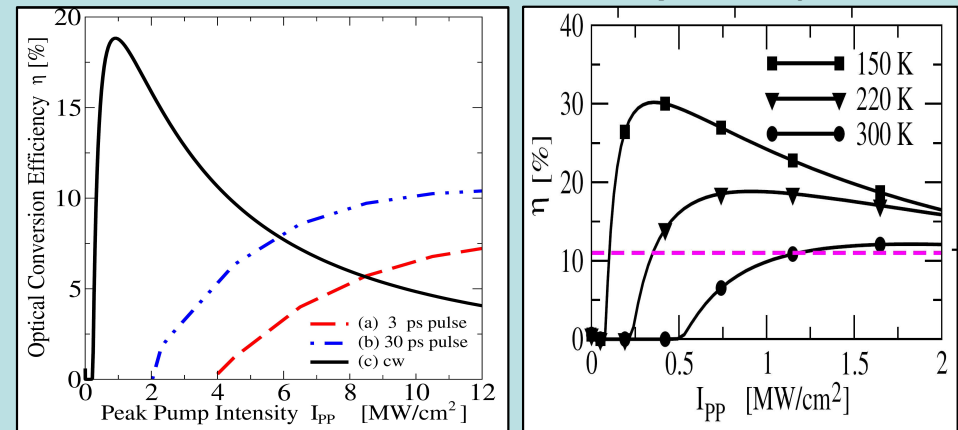
I. OPED-scheme



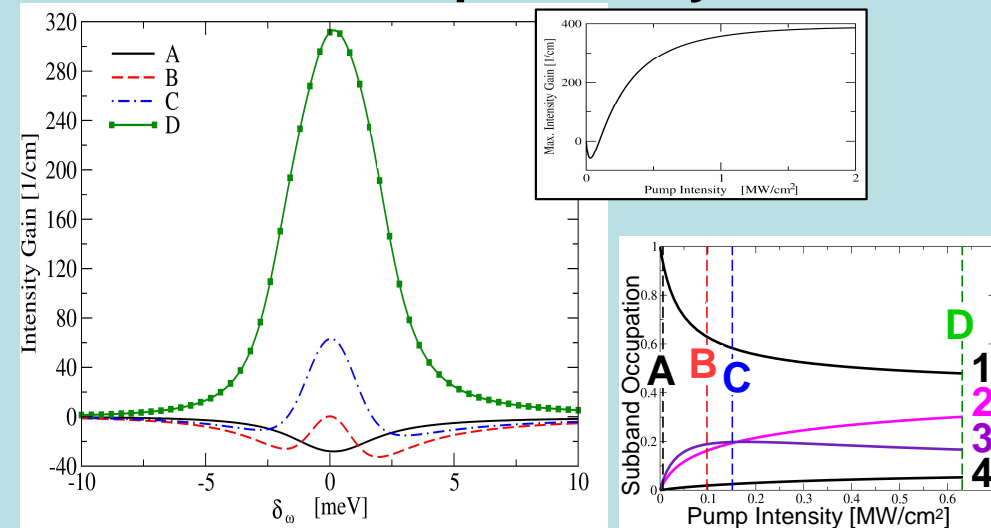
II. Benefits of Pump Recovery



III. Capabilities of OPED (Dependence of Optical Conversion Efficiency on Pulse Duration and Temperature)

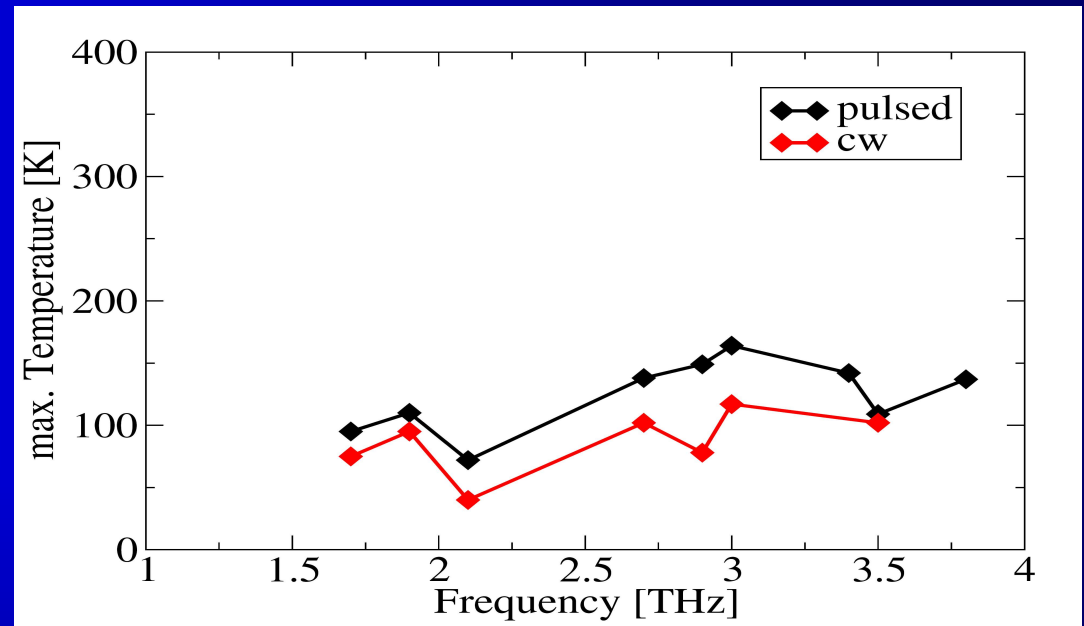
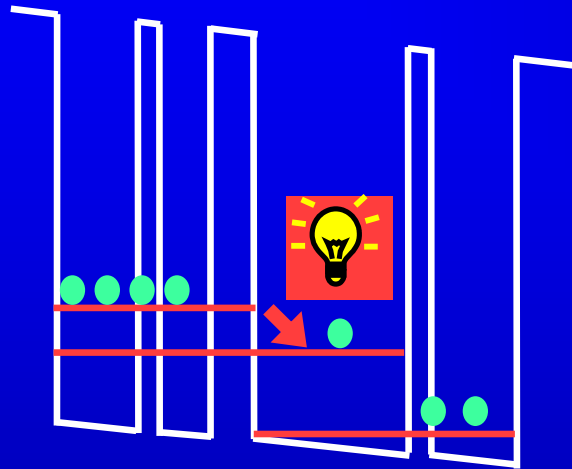


IV. Dependence of Small Signal Gain on Pump Intensity



Motivation: Direct (“conventional”) THz-QCL

direct (“conventional”) THz-QCL



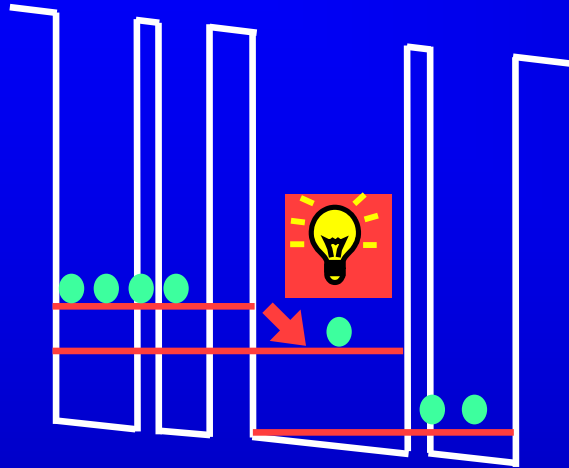
for room temperature (and higher)

hard to achieve/maintain
population inversion

→ efficiency decreases with
increasing temperature

Motivation: Direct (“conventional”) THz-QCL

direct (“conventional”) THz-QCL

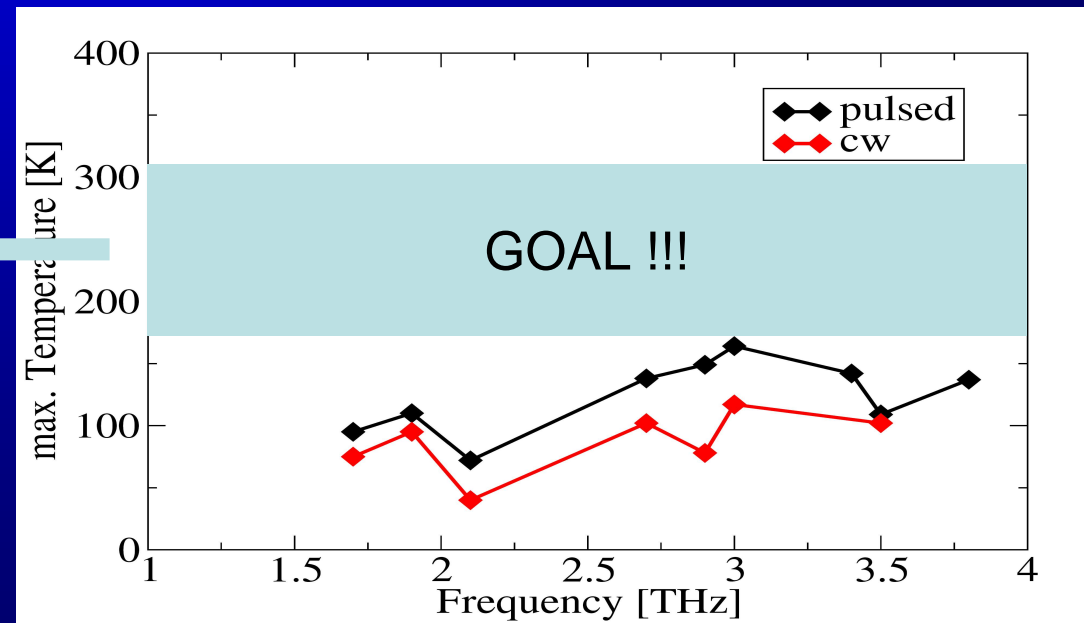


for room temperature (and higher)

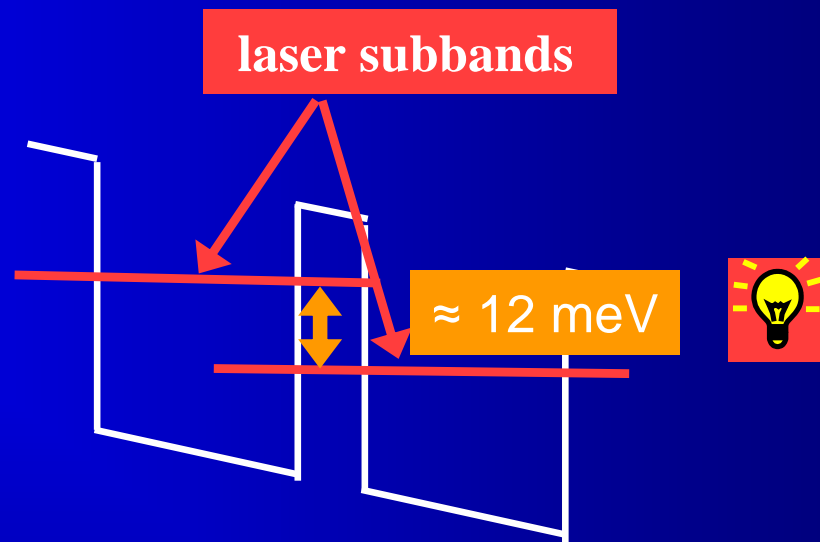
hard to achieve/maintain
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→ efficiency decreases with
increasing temperature

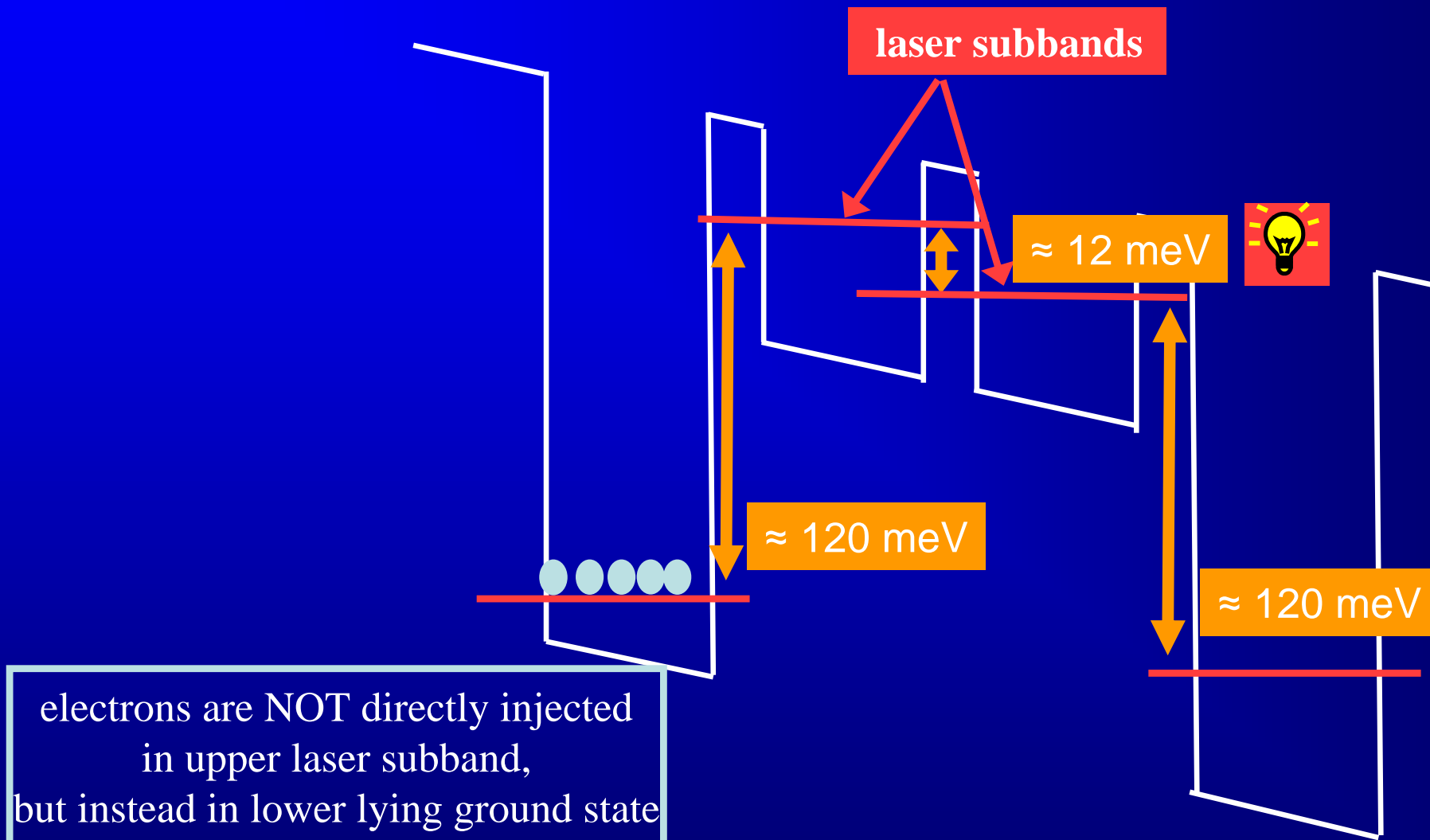
different approach
better suited
for high temperature
operation ?



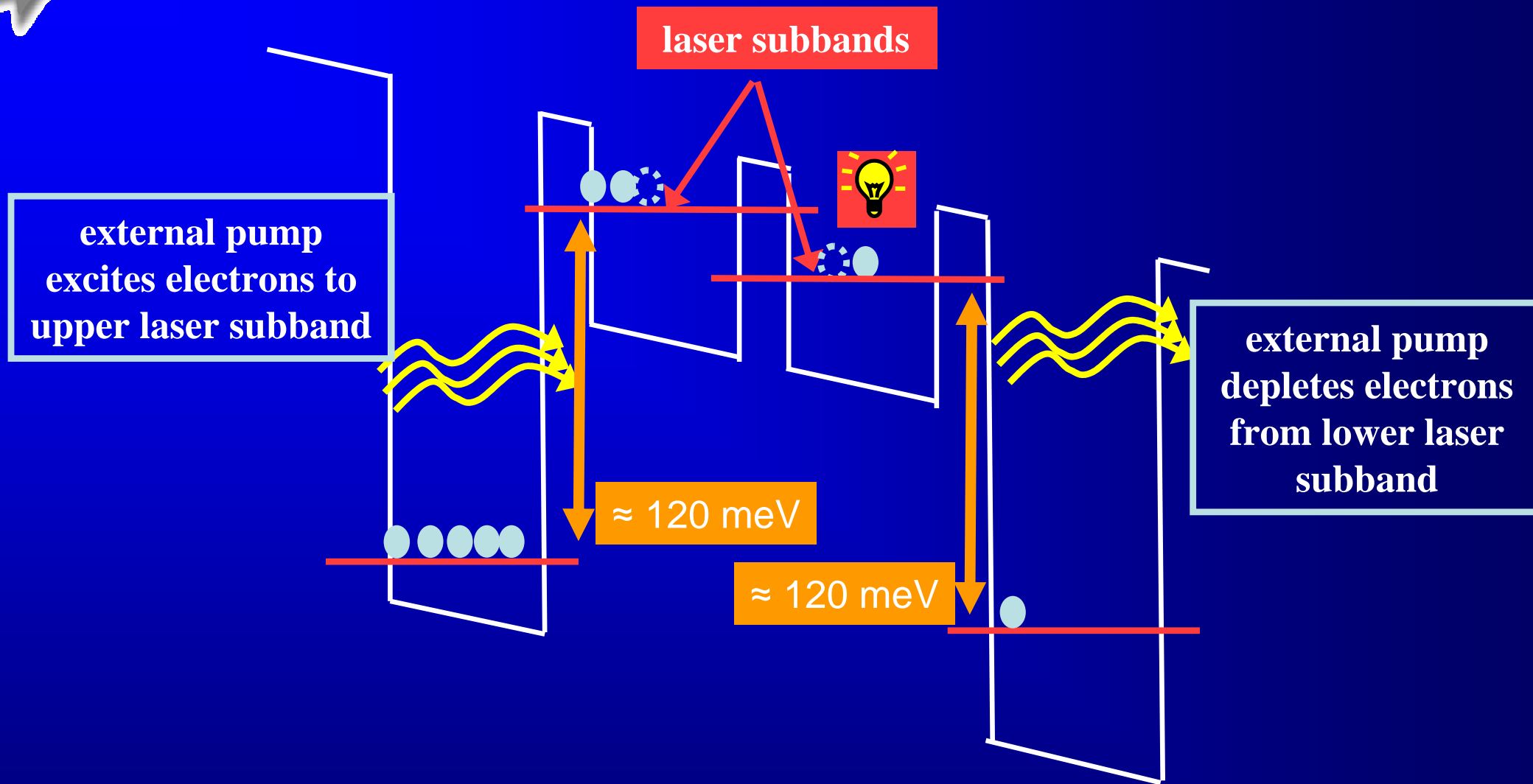
Opt. pumped, electrically driven QCL (OPED)



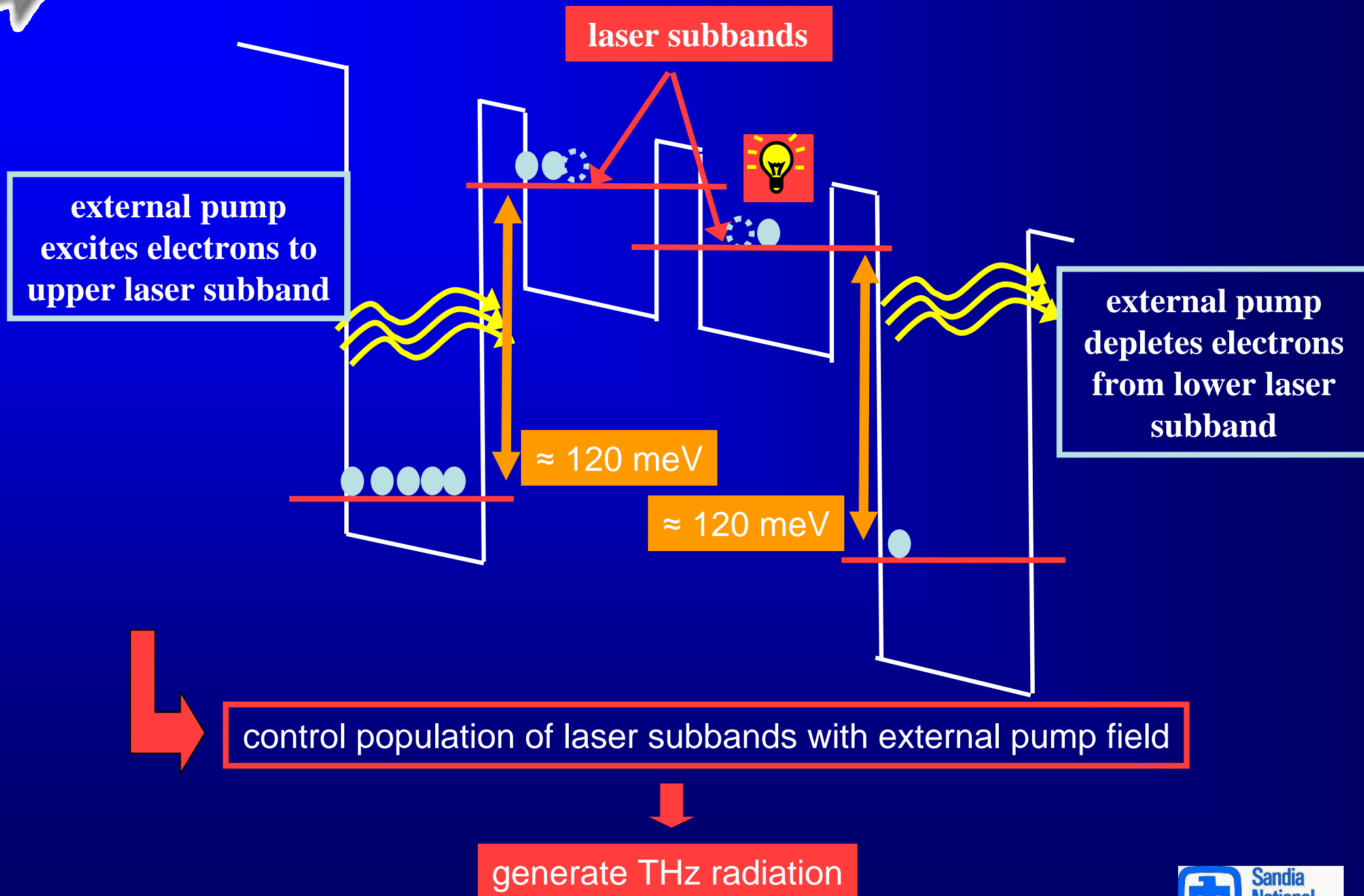
Opt. pumped, electrically driven QCL (OPED)



Opt. pumped, electrically driven QCL (OPED)



Opt. pumped, electrically driven QCL (OPED)



Theory: Maxwell's Equations + SBE

Maxwell's wave equations for total field:

$$\frac{\partial E}{\partial t} = -\frac{1}{\epsilon} \nabla \times H - \frac{1}{\epsilon} \frac{\partial P}{\partial t}, \quad \frac{\partial H}{\partial t} = -\frac{1}{\mu} \nabla \times E$$

E, H

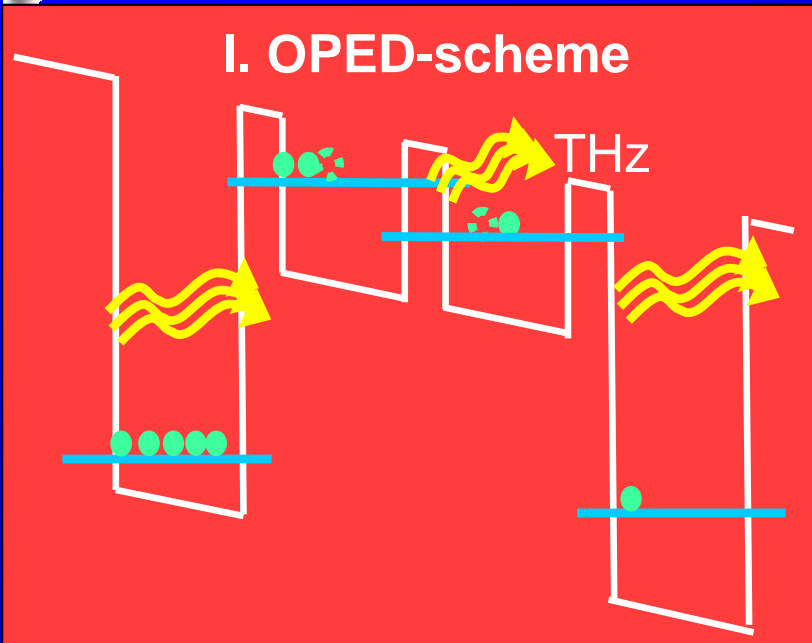
$$P = \sum_{ij} d_{ij} P_{ij}$$

SBE for polarizations P and subband populations N :

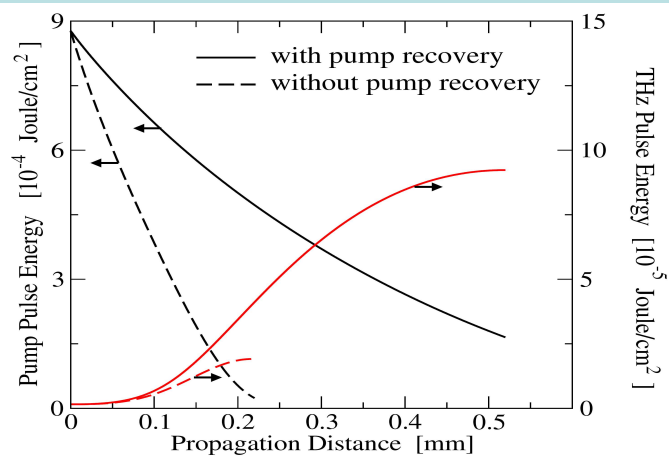
$$\begin{aligned} \frac{d}{dt} P_{ab} &= \left(\frac{i}{\hbar} \epsilon_{ab} - \gamma \right) P_{ab} + \frac{i}{\hbar} d_{ab} (N_a - N_b) E + \frac{i}{\hbar} \sum_{c \neq a, b} (d_{bc} P_{ac} - d_{ac} P_{cb}) E \\ \frac{d}{dt} N_a &= \frac{i}{2\hbar} \sum_{c \neq a} d_{ac} (P_{ca} - P_{ac}) E + \frac{d}{dt} N_a|_s \end{aligned}$$

Overview

I. OPED-scheme

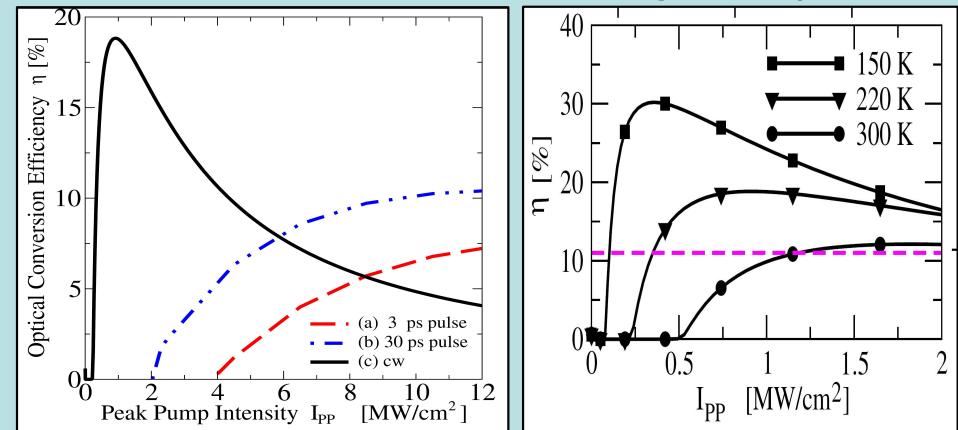


II. Benefits of Pump Recovery

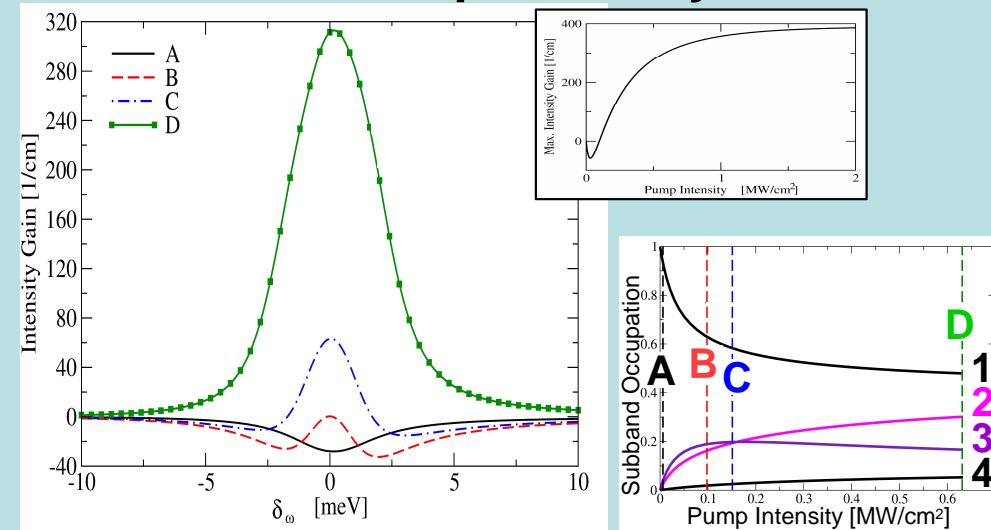


III. Capabilities of OPED

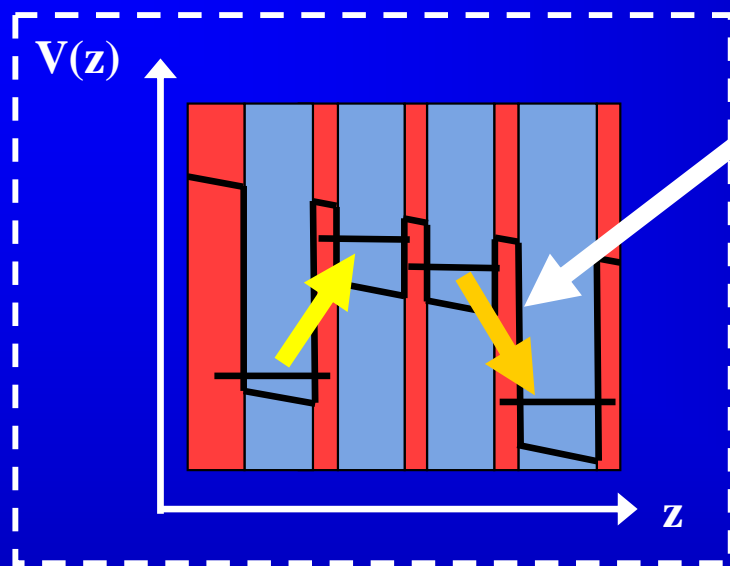
(Dependence of Optical Conversion Efficiency on Pulse Duration and Temperature)



IV. Dependence of Small Signal Gain on Pump Intensity

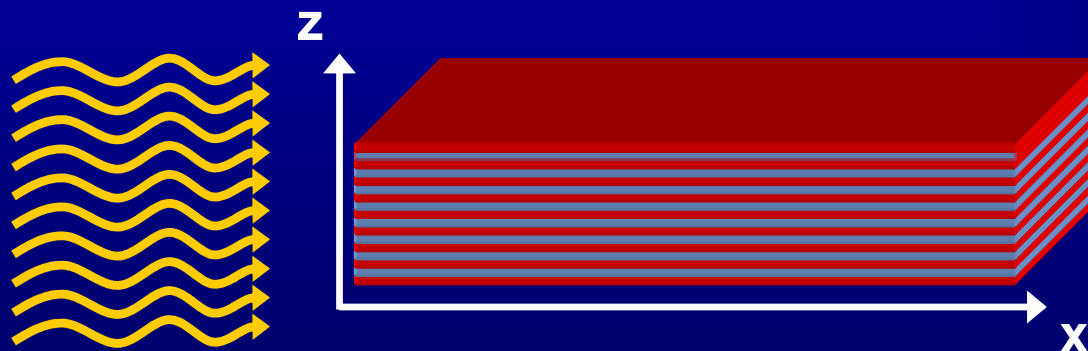


Benefits of Pump Recovery

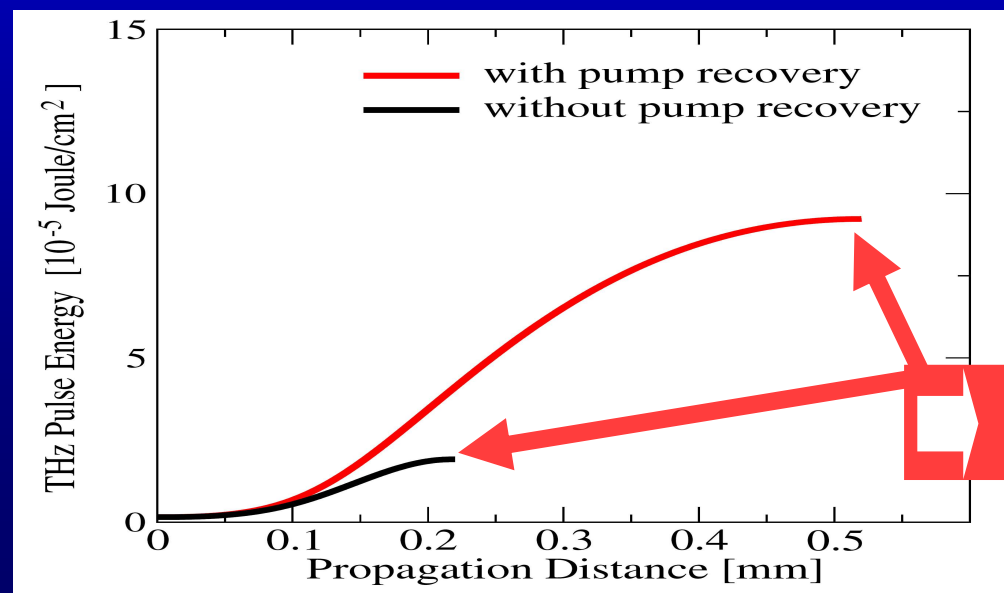
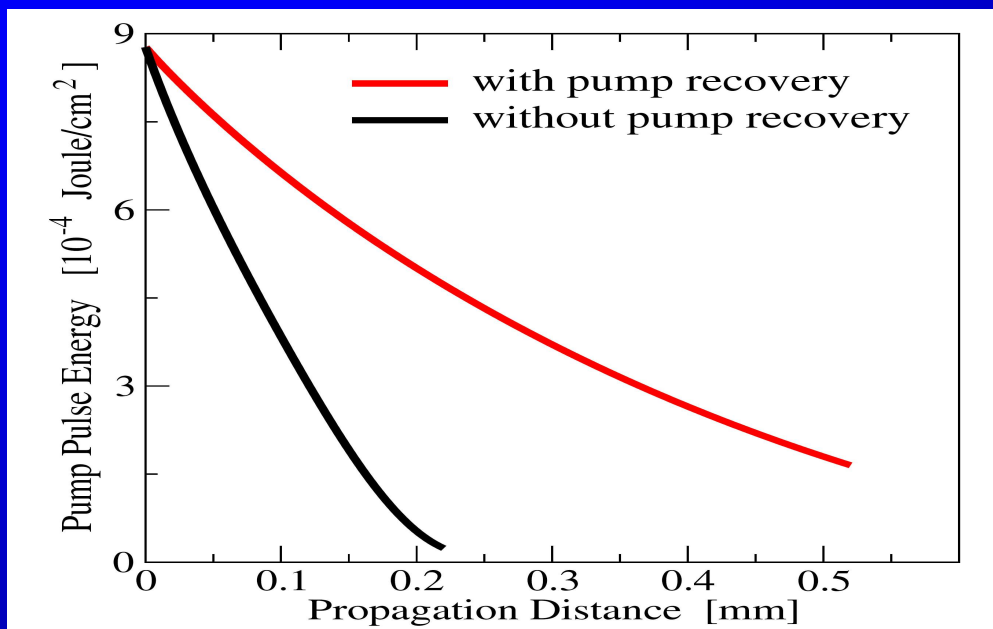


Coherent Pump Recovery !!!

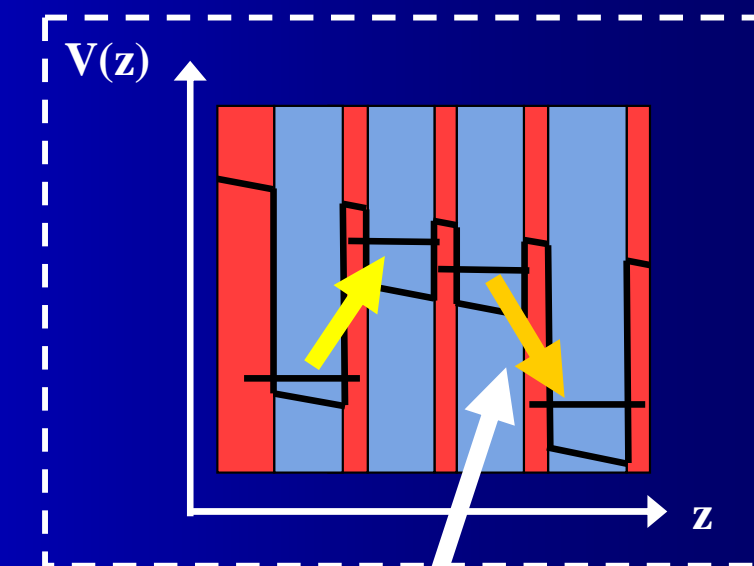
excite structure with Gaussian pulse (propagating in x-direction)



Benefits of Pump Recovery



4-5 x more THz radiation

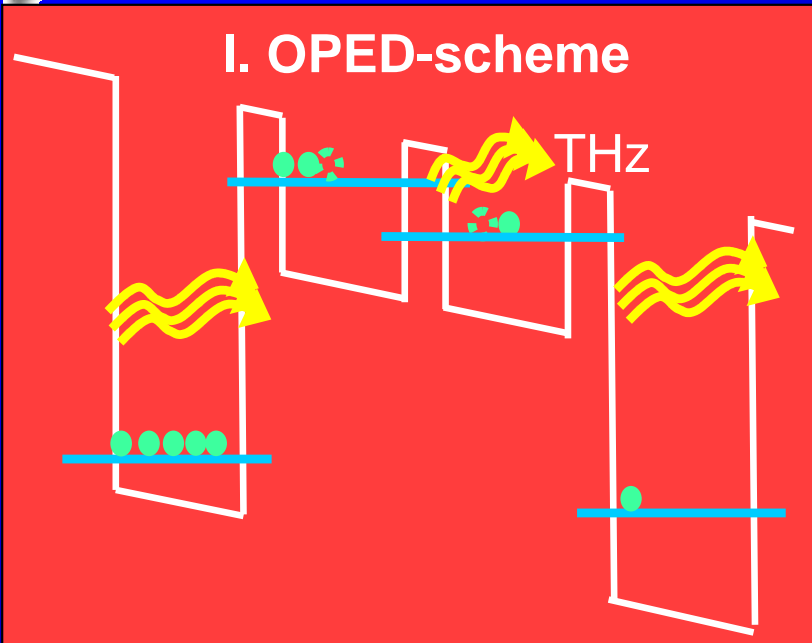


Coherent Pump Recovery !!!

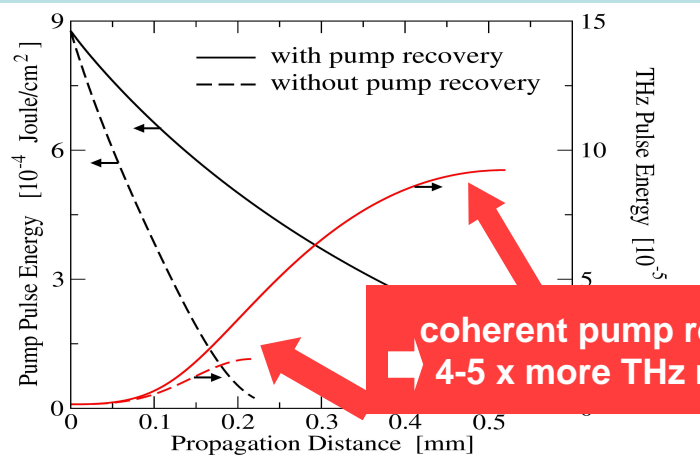
total losses for simulations:
 $\alpha_{THz} = 140cm^{-1}$, $\alpha_{FIR} = 50cm^{-1}$

Overview

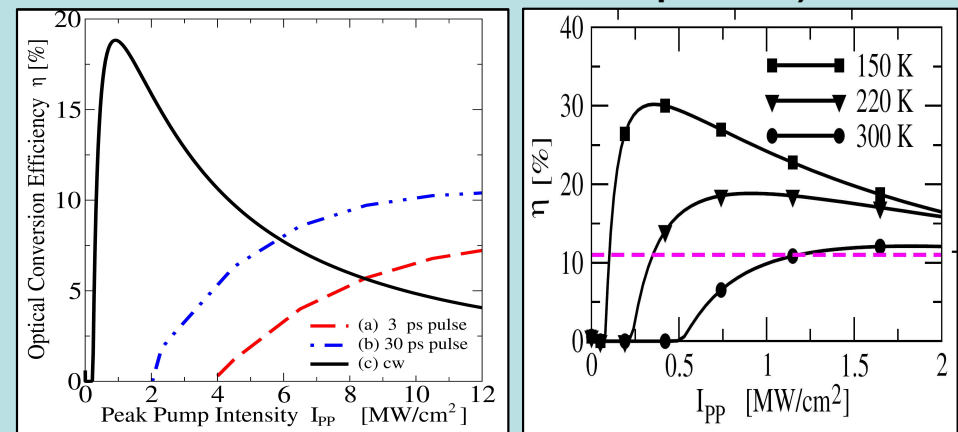
I. OPED-scheme



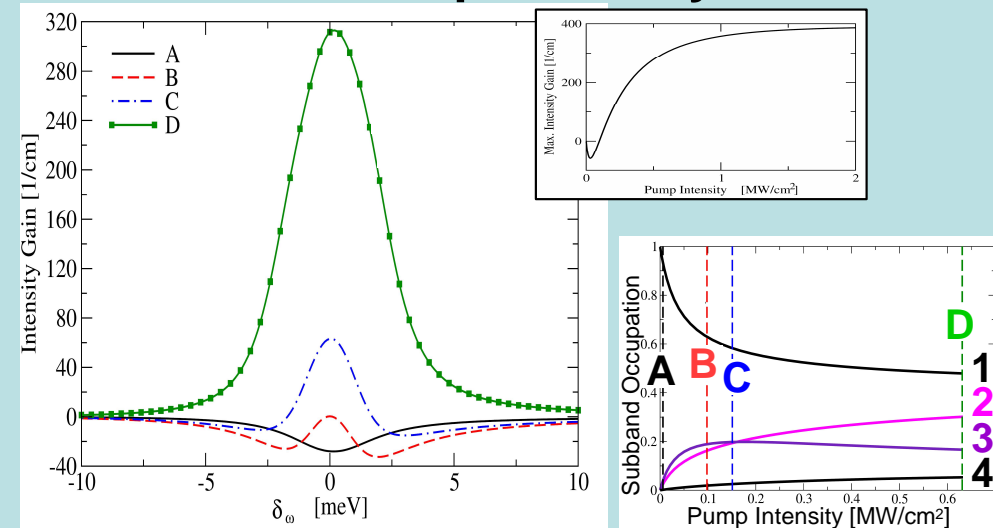
II. Benefits of Pump Recovery



III. Capabilities of OPED (Dependence of Optical Conversion Efficiency on Pulse Duration and Temperature)



IV. Dependence of Small Signal Gain on Pump Intensity



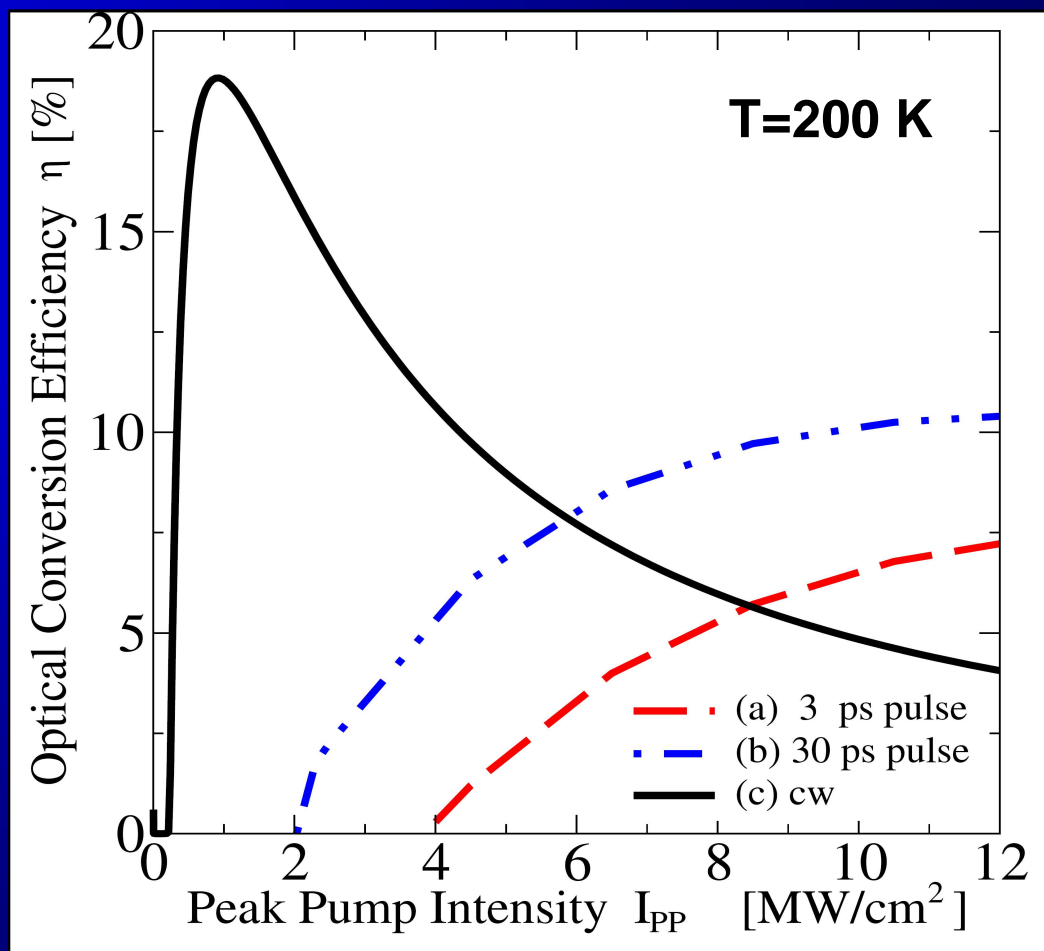
Optical Conversion Efficiency (max. Output/Input)

Dependence of optical conversion efficiency η on pump energy

for (a) excitation with 3 ps pulse
for (b) excitation with 30 ps pulse
for (c) cw excitation

- strong dependence on excitation length and intensity
- saturation of THz gain (optical Stark effect and pump-induced population redistribution)

window for optimal pump excitation



total losses for simulations:

$$\alpha_{THz} = 140 \text{ cm}^{-1}, \alpha_{FIR} = 50 \text{ cm}^{-1}$$

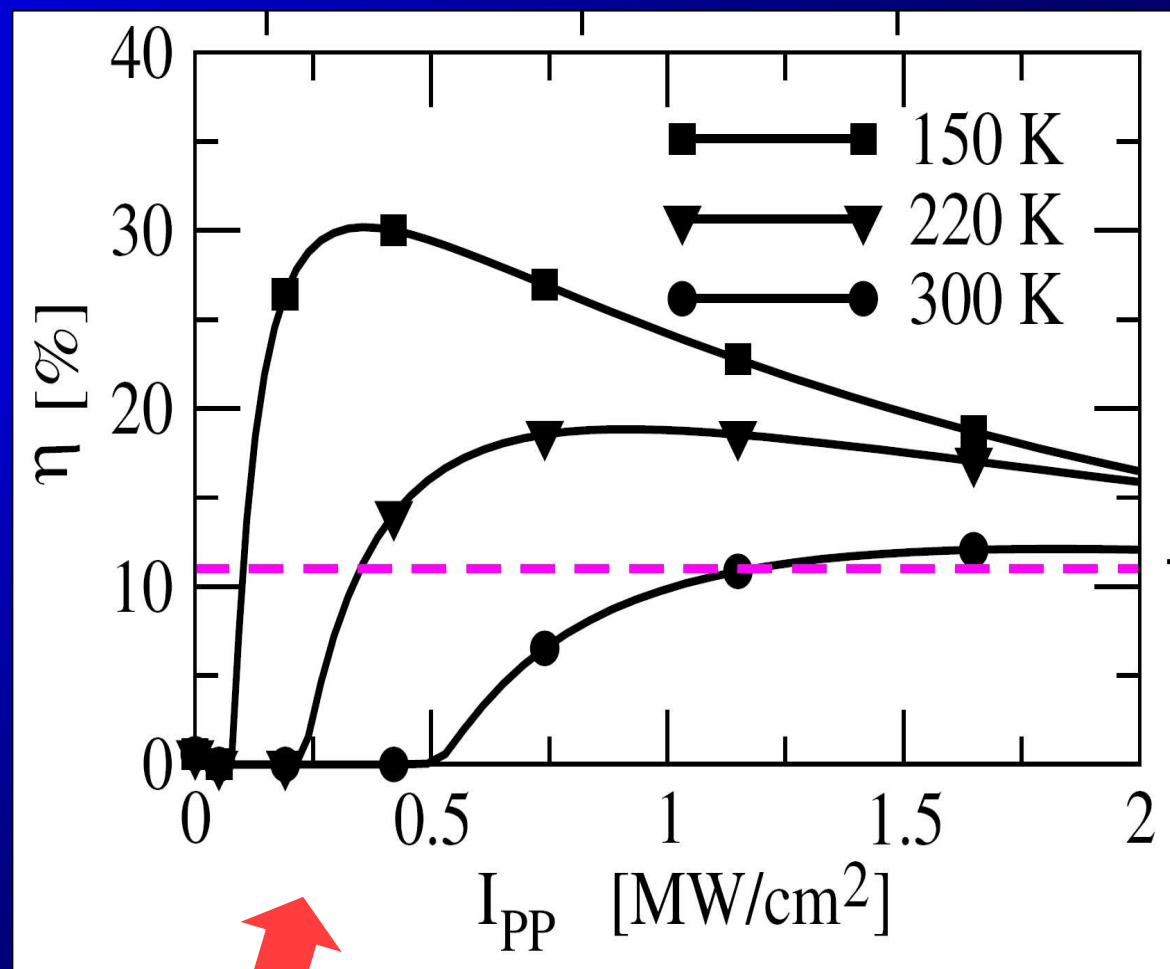
Optical Conversion Efficiency

Dependence of optical conversion efficiency η on pump energy

for cw excitation
at different temperatures

strong temp. dependence due to
temp. dependent non-radiative
carrier recombination and
polarization dephasing

highest achievable η and
width of the window
for optimal pump excitation
are temperature dependent

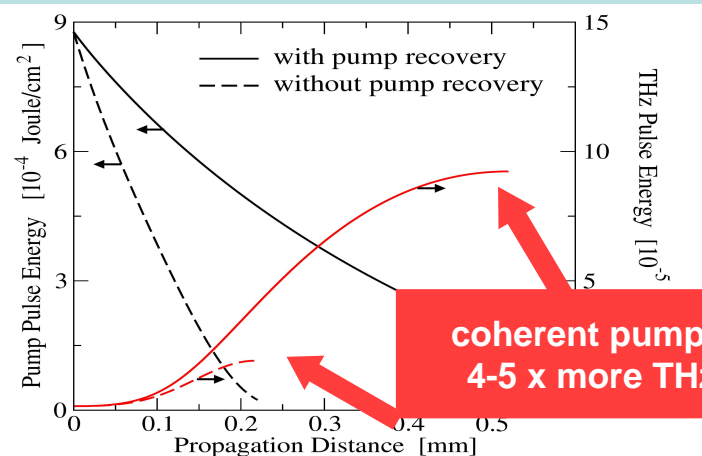


**cw-excitation:
pathway to room-temp operation**

total losses for simulations:
 $\alpha_{THz} = 140\text{cm}^{-1}$, $\alpha_{FIR} = 50\text{cm}^{-1}$

Promising Results for Pulse- and CW-Excitation:

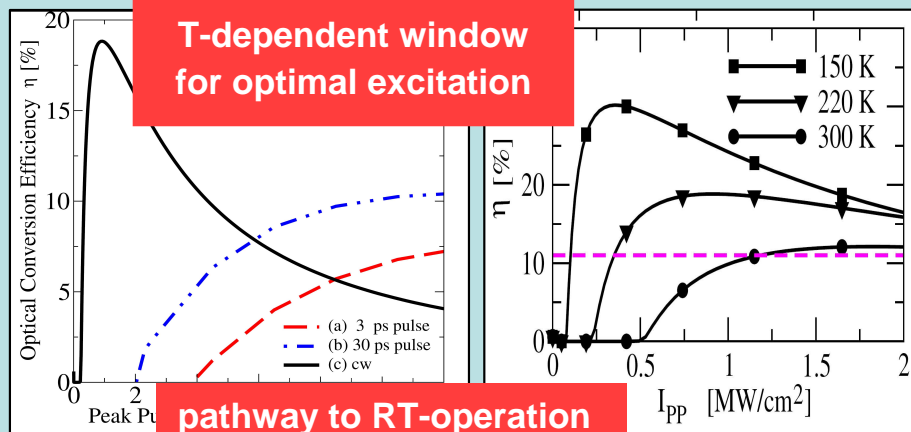
II. Benefits of Pump Recovery



good recovery of pump energy

III. Capabilities of OPED

(Dependence of Optical Conversion Efficiency on Pulse Duration and Temperature)



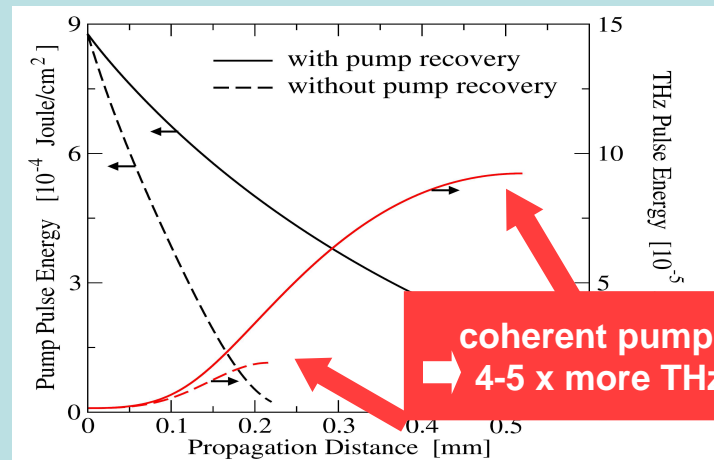
promising conversion efficiencies for different lattice temperatures

total losses for simulations:

$$\alpha_{THz} = 140 cm^{-1}, \alpha_{FIR} = 50 cm^{-1}$$

Promising Results for Pulse- and CW-Excitation:

II. Benefits of Pump Recovery



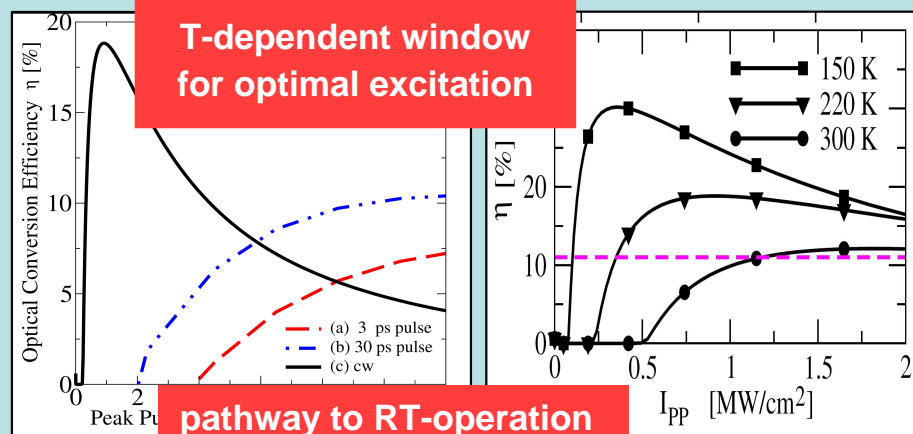
good recovery of pump energy

This means that our approach works -
- theoretically -
like we thought it would!
Or does it?

promising conversion efficiencies
for different lattice temperatures

III. Capabilities of OPED

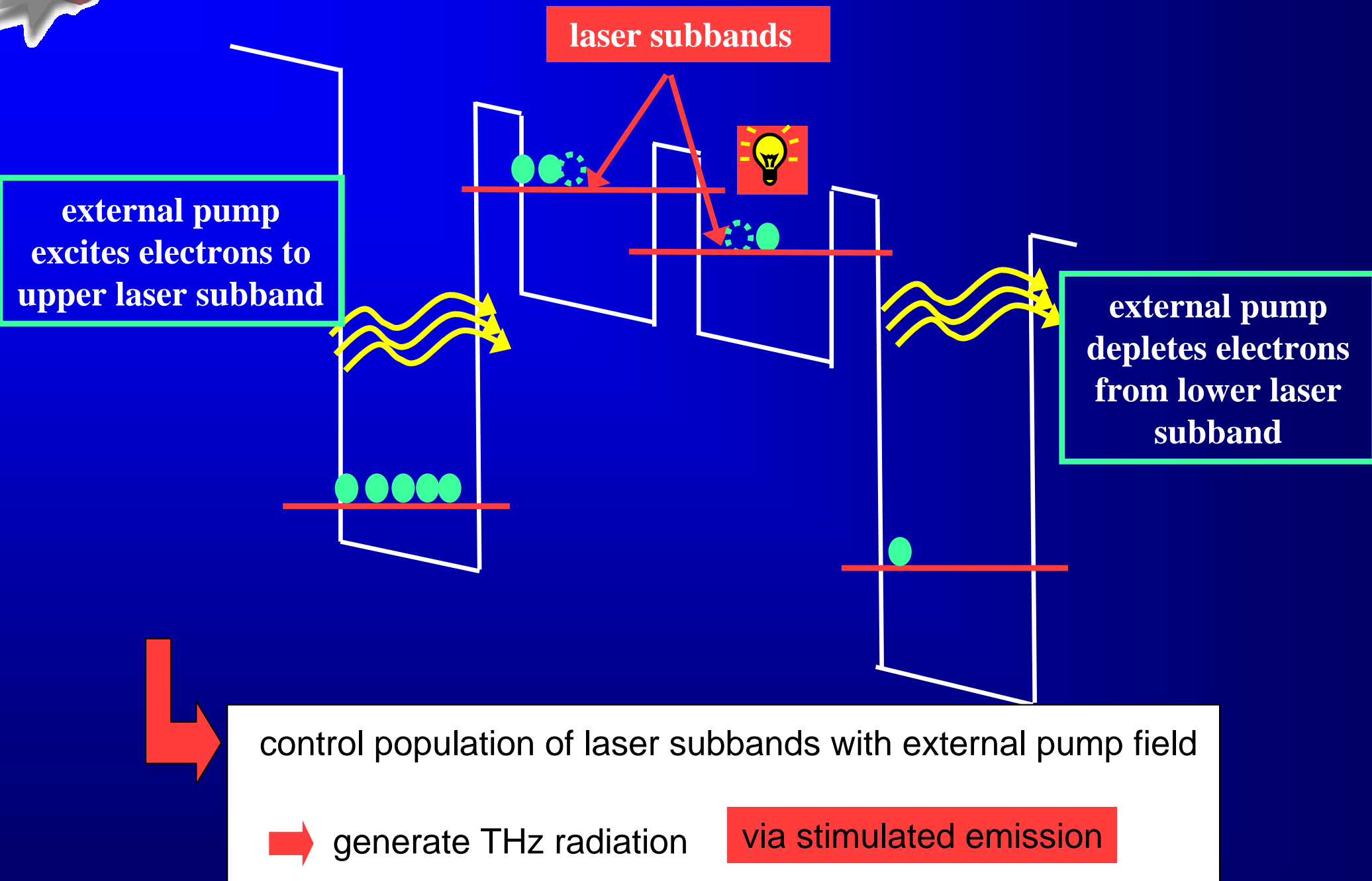
(Dependence of Optical Conversion Efficiency on Pulse Duration and Temperature)



total losses for simulations:

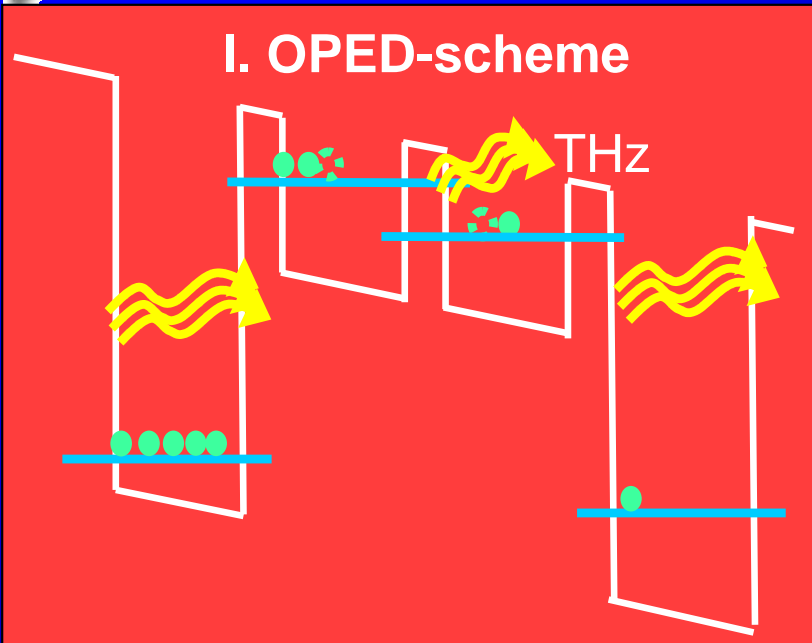
$$\alpha_{THz} = 140 cm^{-1}, \alpha_{FIR} = 50 cm^{-1}$$

It works like we thought it would! Or does it?

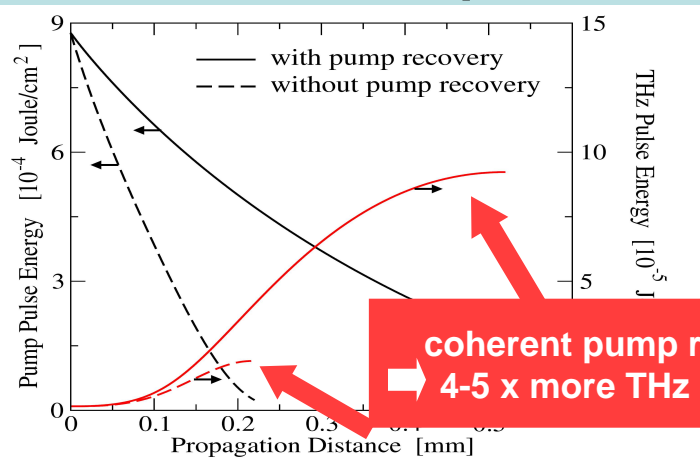


Overview

I. OPED-scheme

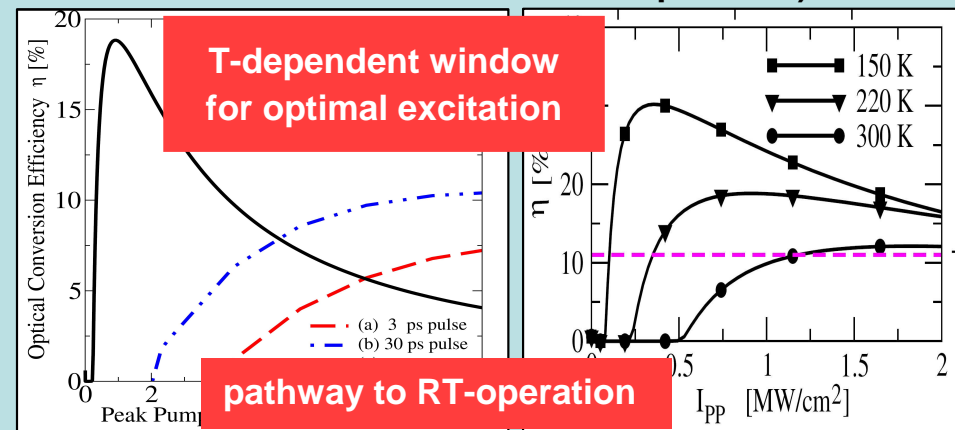


II. Benefits of Pump Recovery

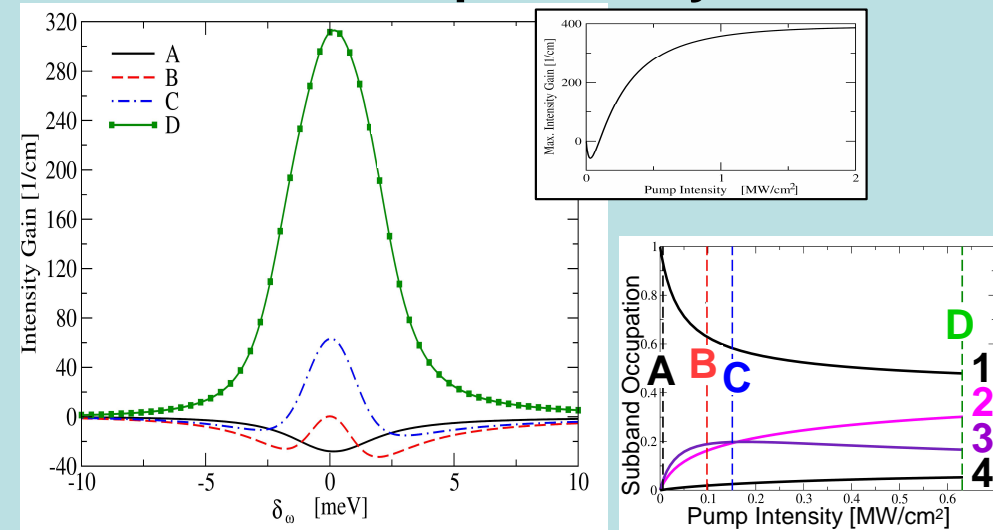


III. Capabilities of OPED

(Dependence of Optical Conversion Efficiency on Pulse Duration and Temperature)

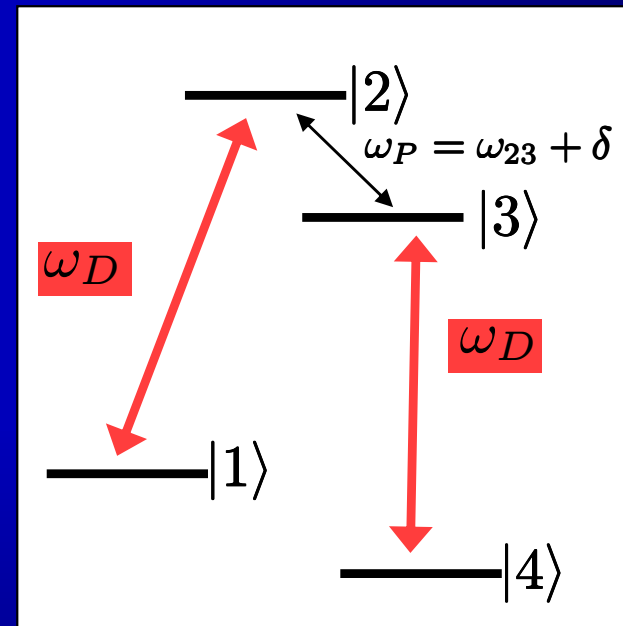
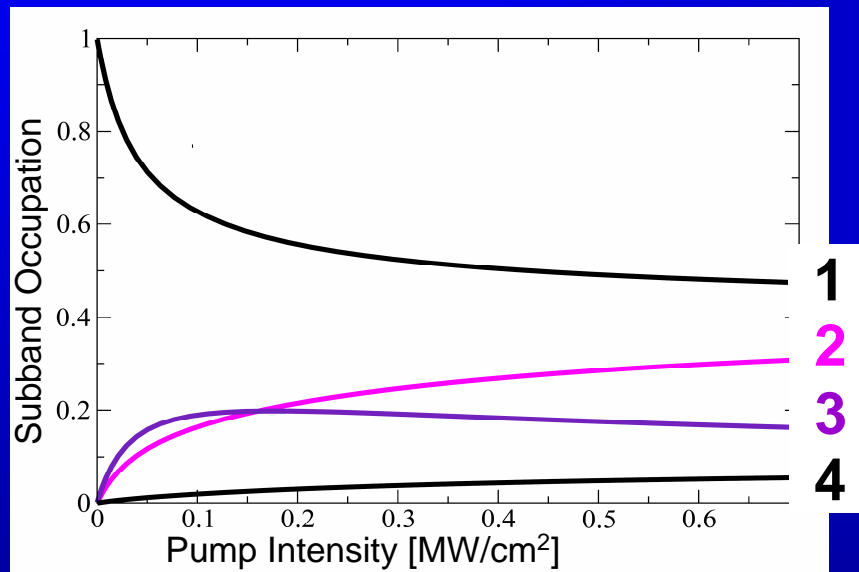


IV. Dependence of Small Signal Gain on Pump Intensity



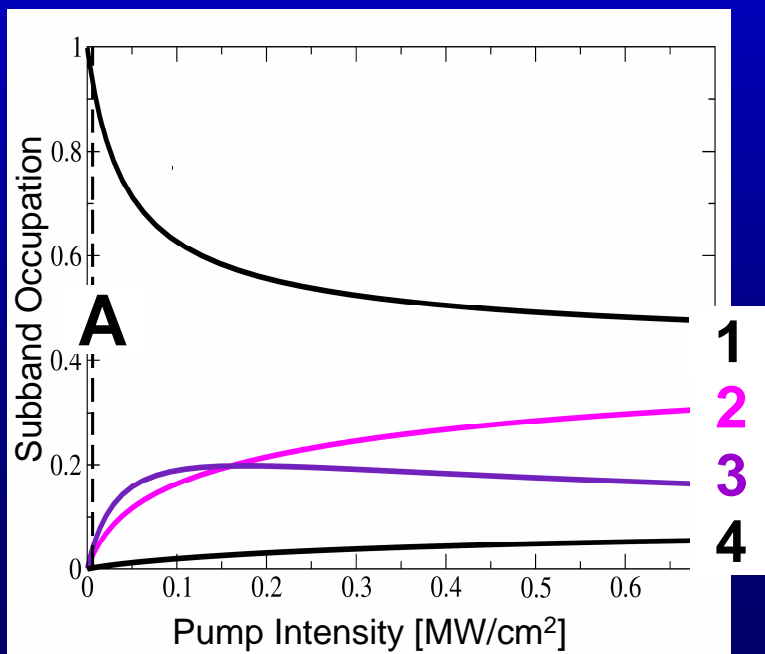
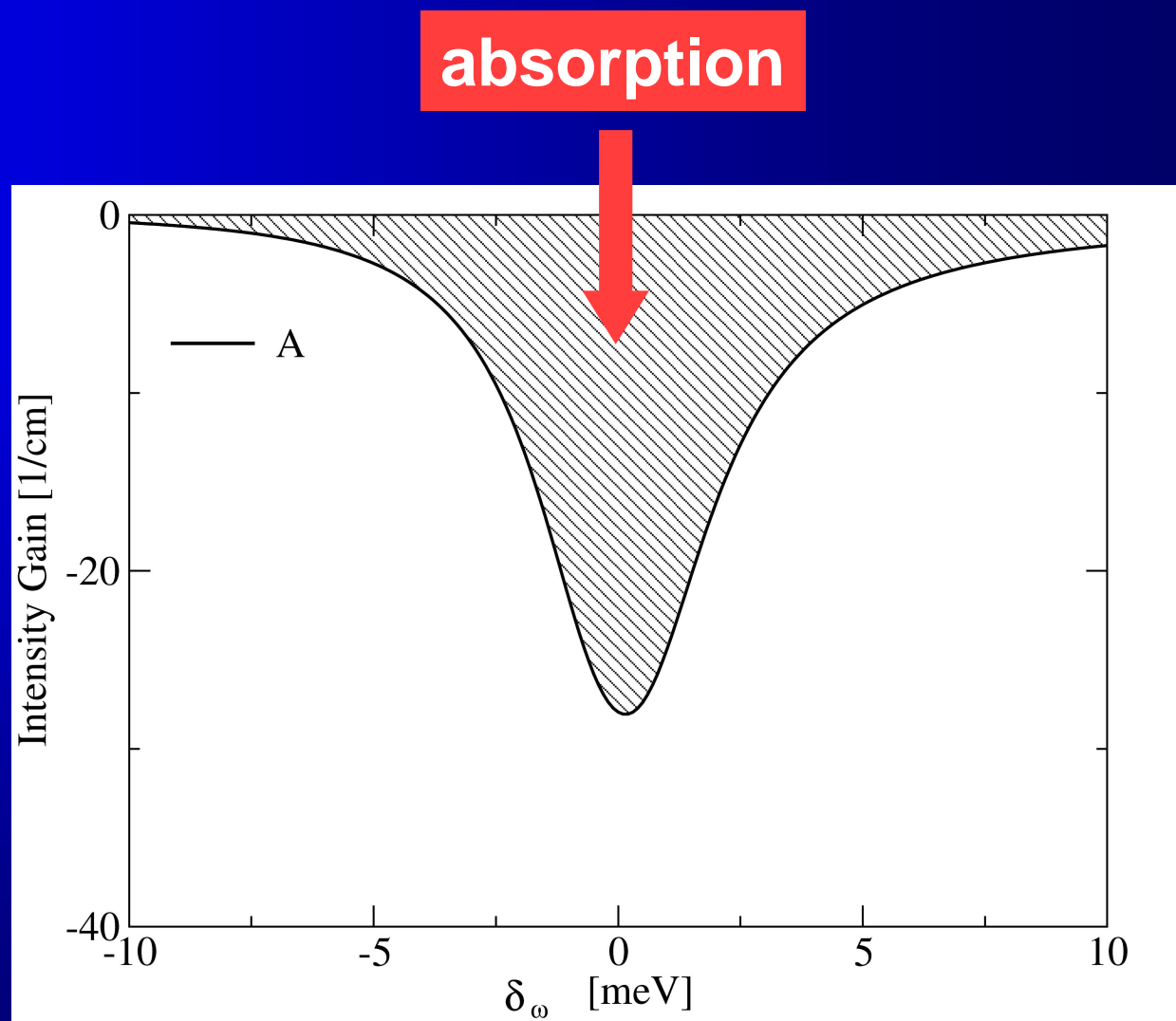
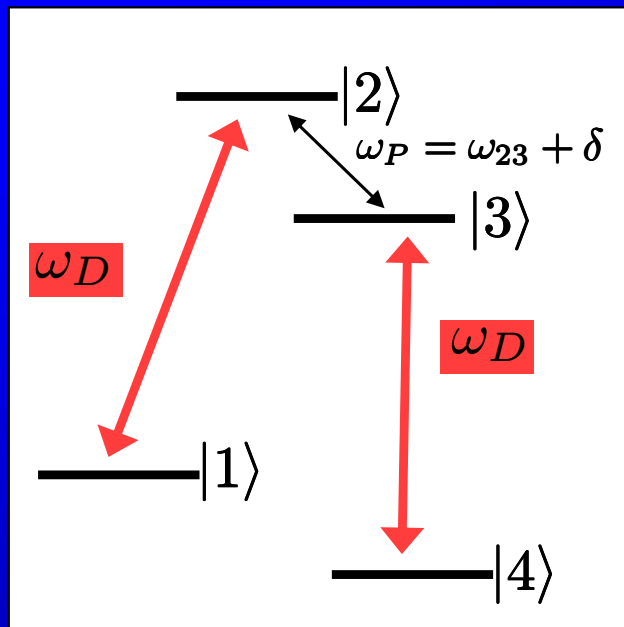
Let's look at the linear gain for different pump intensities

- excite OPED with pump field
(redistribute carriers between subbands)

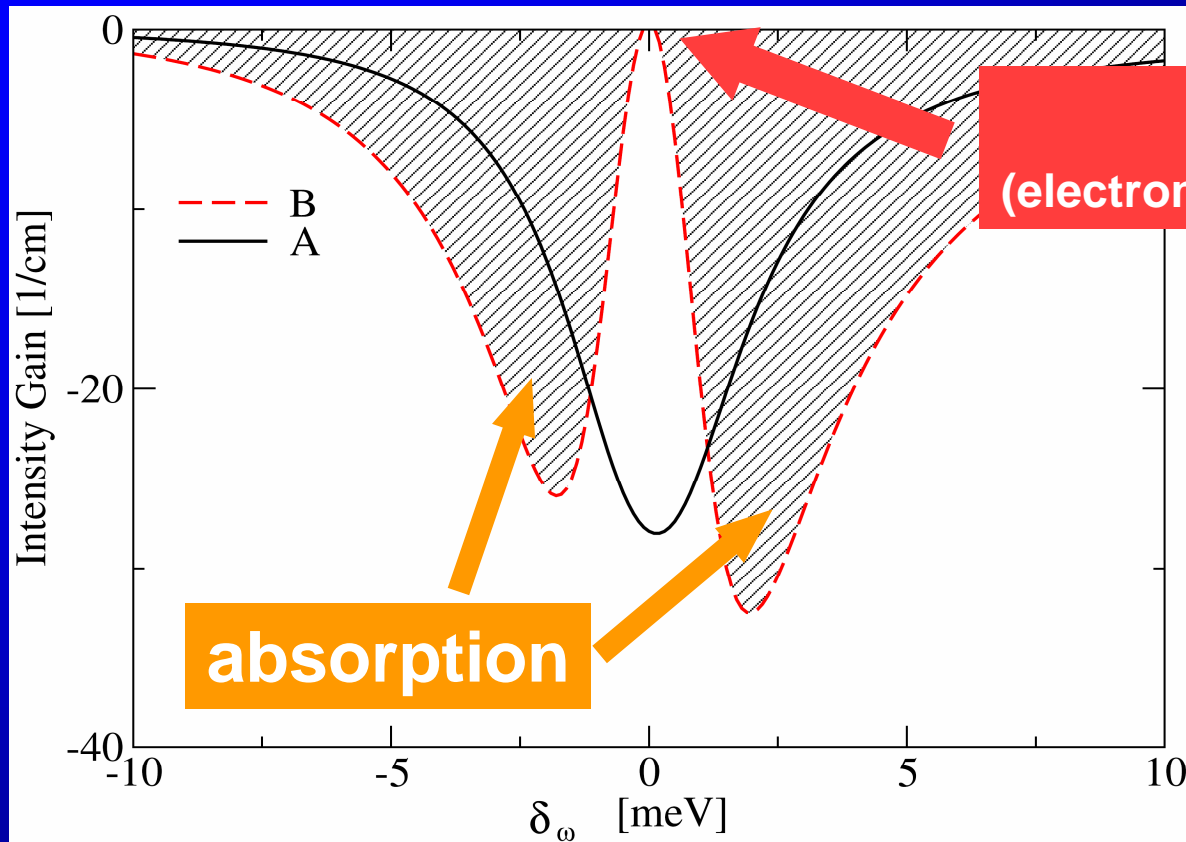


- probe with small THz field with different detunings

Let's look at the linear gain for different drive intensities

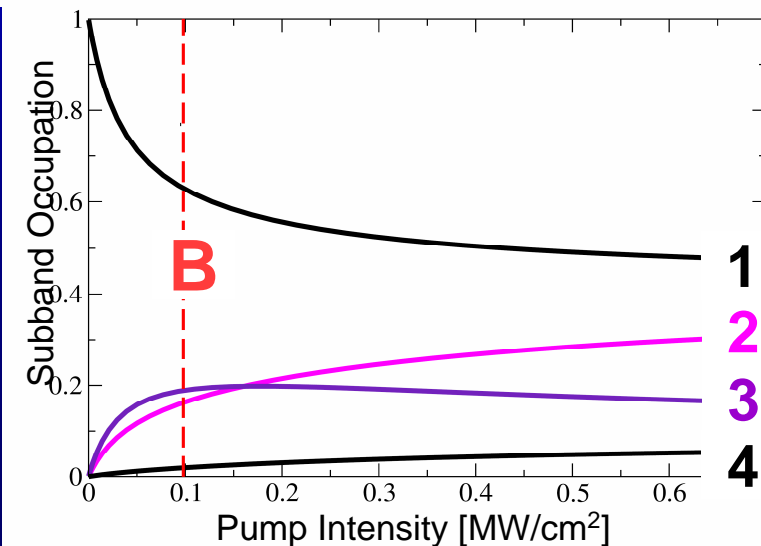


Let's look at the linear gain for different drive intensities

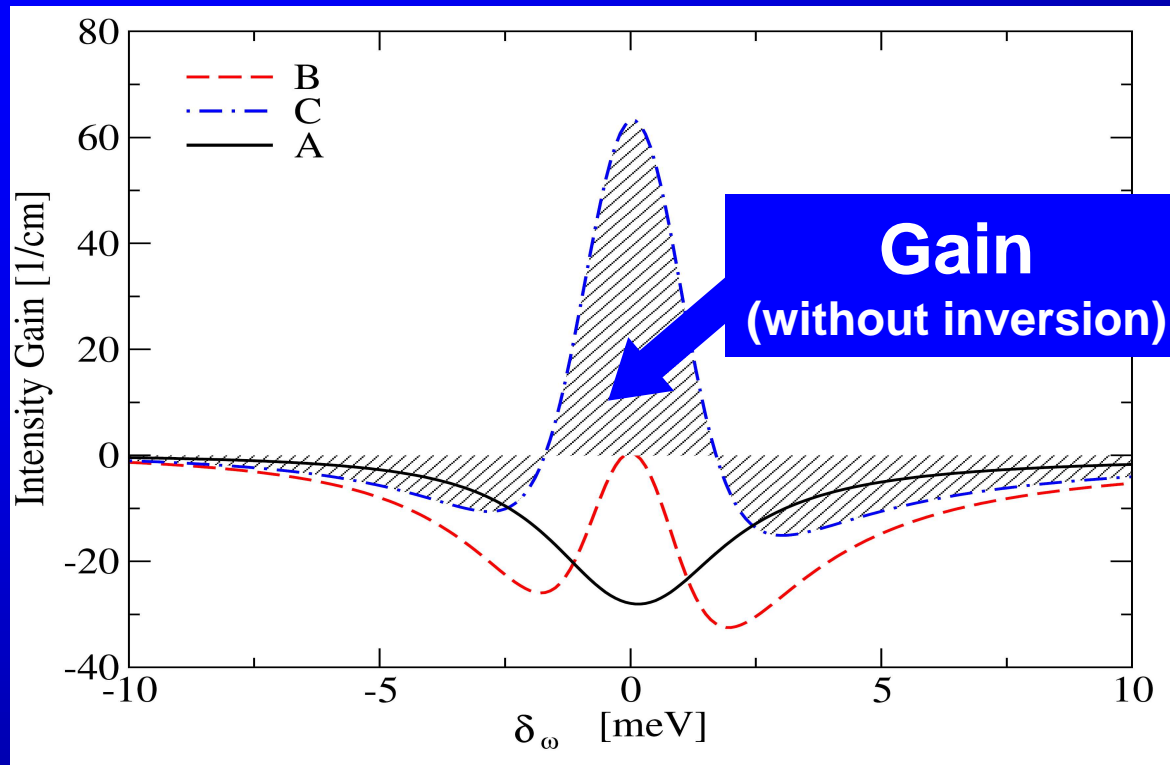


EIT
(electromagnetically induced transparency)

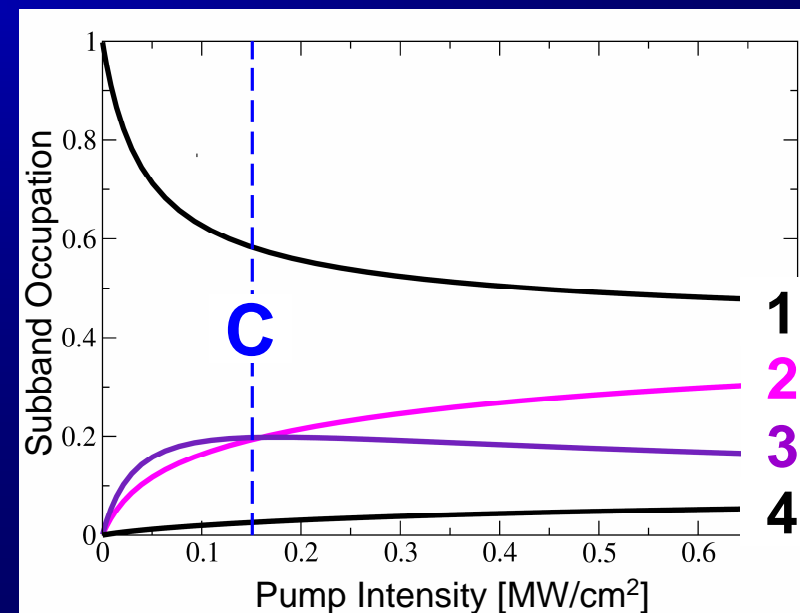
SIGN OF QUANTUM COHERENCE!!!



Let's look at the linear gain for different drive intensities

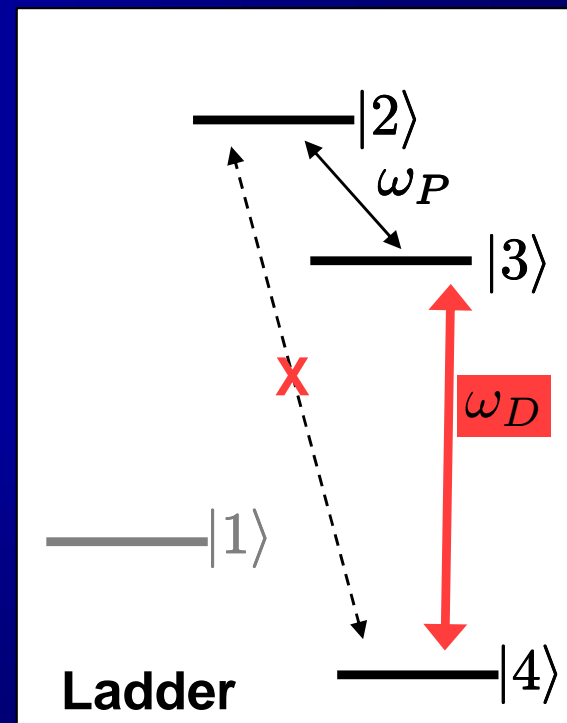
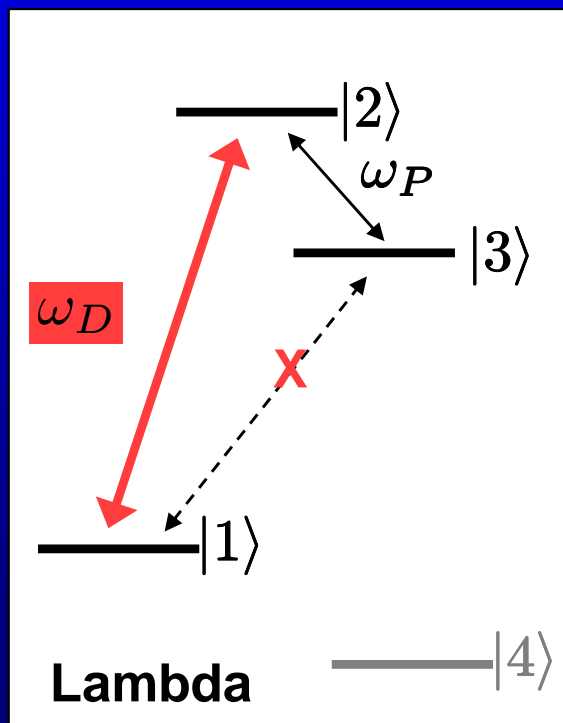


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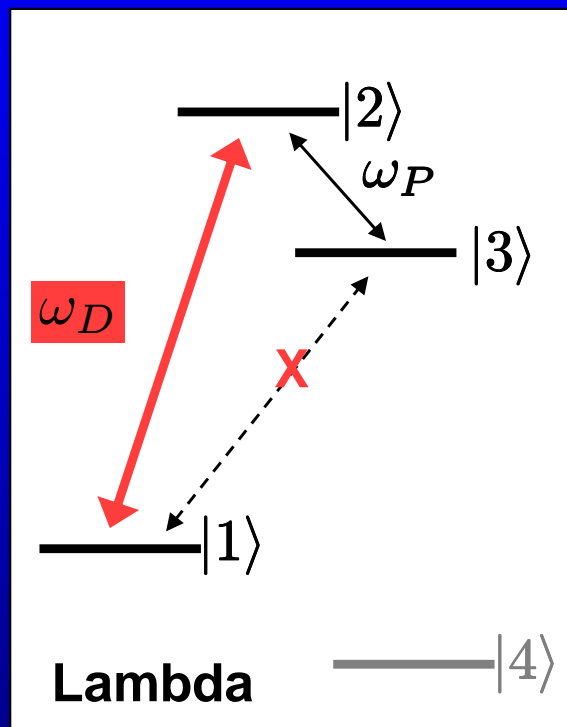


Quantum Coherence Contributions

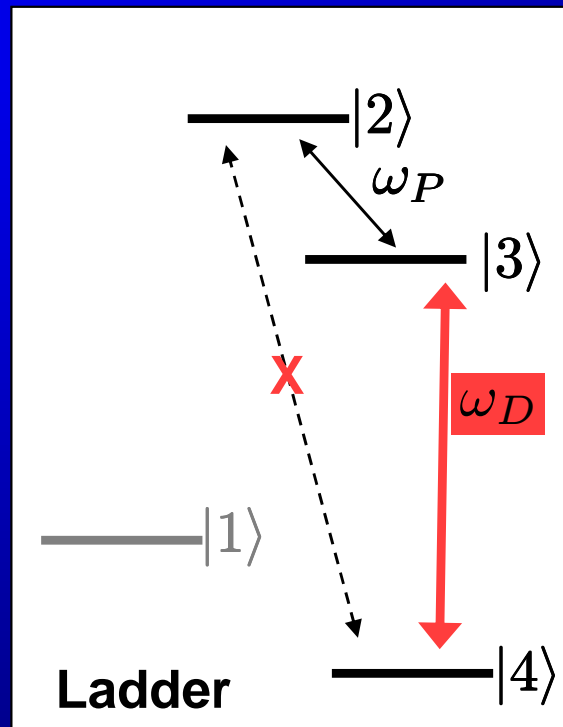
Quantum coherence contributions are usually observed in atomic model system, such as ... e.g. lambda and ladder systems



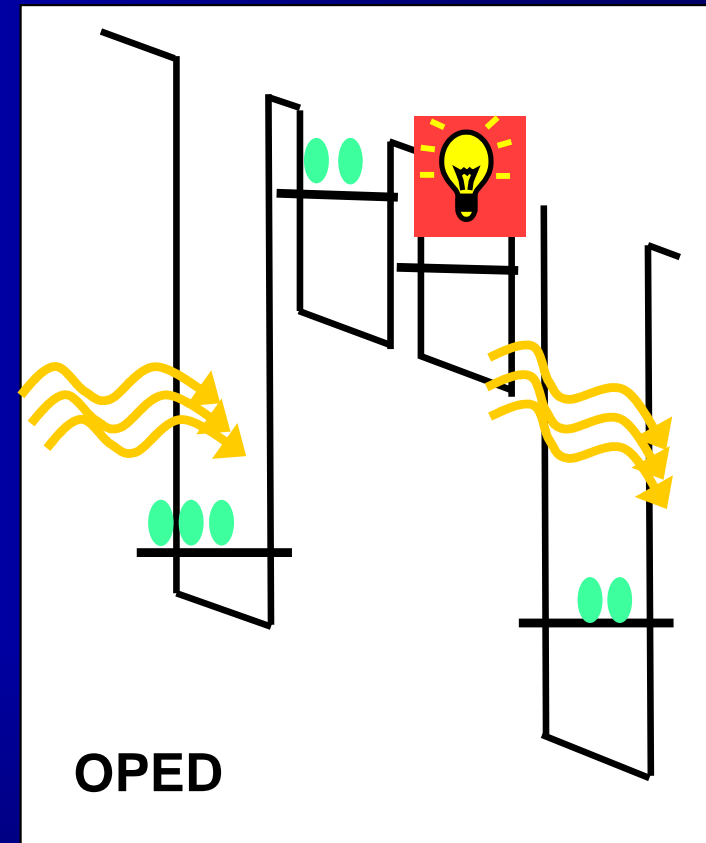
Lambda + Ladder = OPED



+



=



model systems for quantum coherence effects
(LWI, EIT)
should see signs of q. c. in gain spectra

Analytical Results:

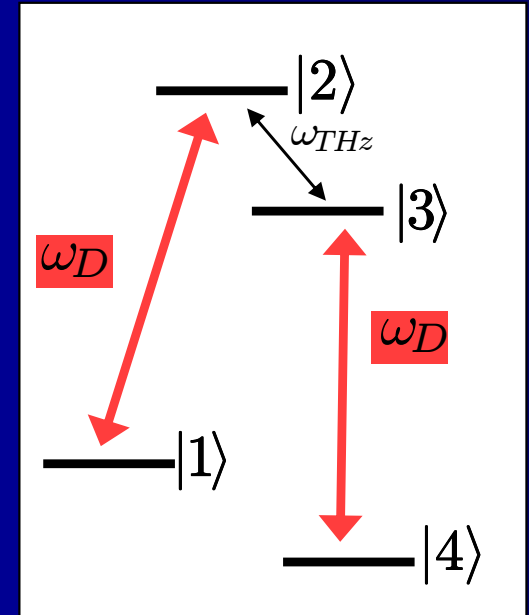
- possible to generate THz radiation with/without population inversion due to:
- stimulated emission
 - quantum coherence contributions

linear stimulated emission

automatically phase-matched quantum coherence effects

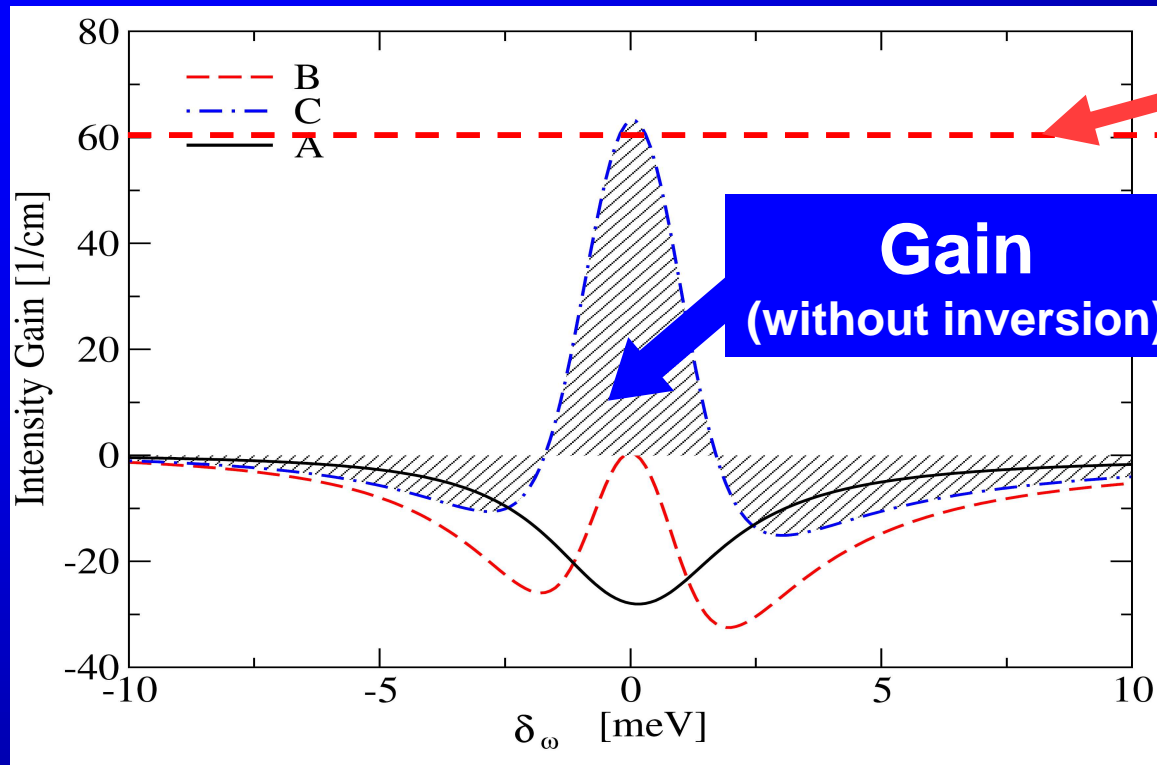
auto-correlation

cross-correlation



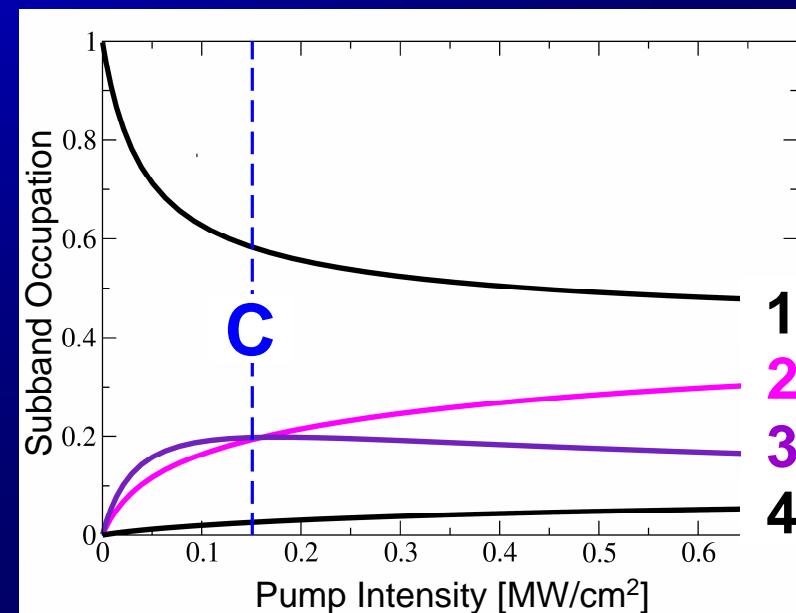
$$\frac{d}{dt} E_{THz} \propto d_{THz}^2 E_{THz} \frac{(N_2 - N_3)4\hbar^2\gamma^2 + (N_2 - N_3)d_{THz}^2|E_{THz}|^2 + (N_1 + N_2 - N_3 - N_4)d_D^2|E_D|^2}{2\hbar\gamma(4d_D^2|E_D|^2 + d_{THz}^2|E_{THz}|^2 + 4\hbar^2\gamma^2)}$$

Let's look at the linear gain for different drive intensities

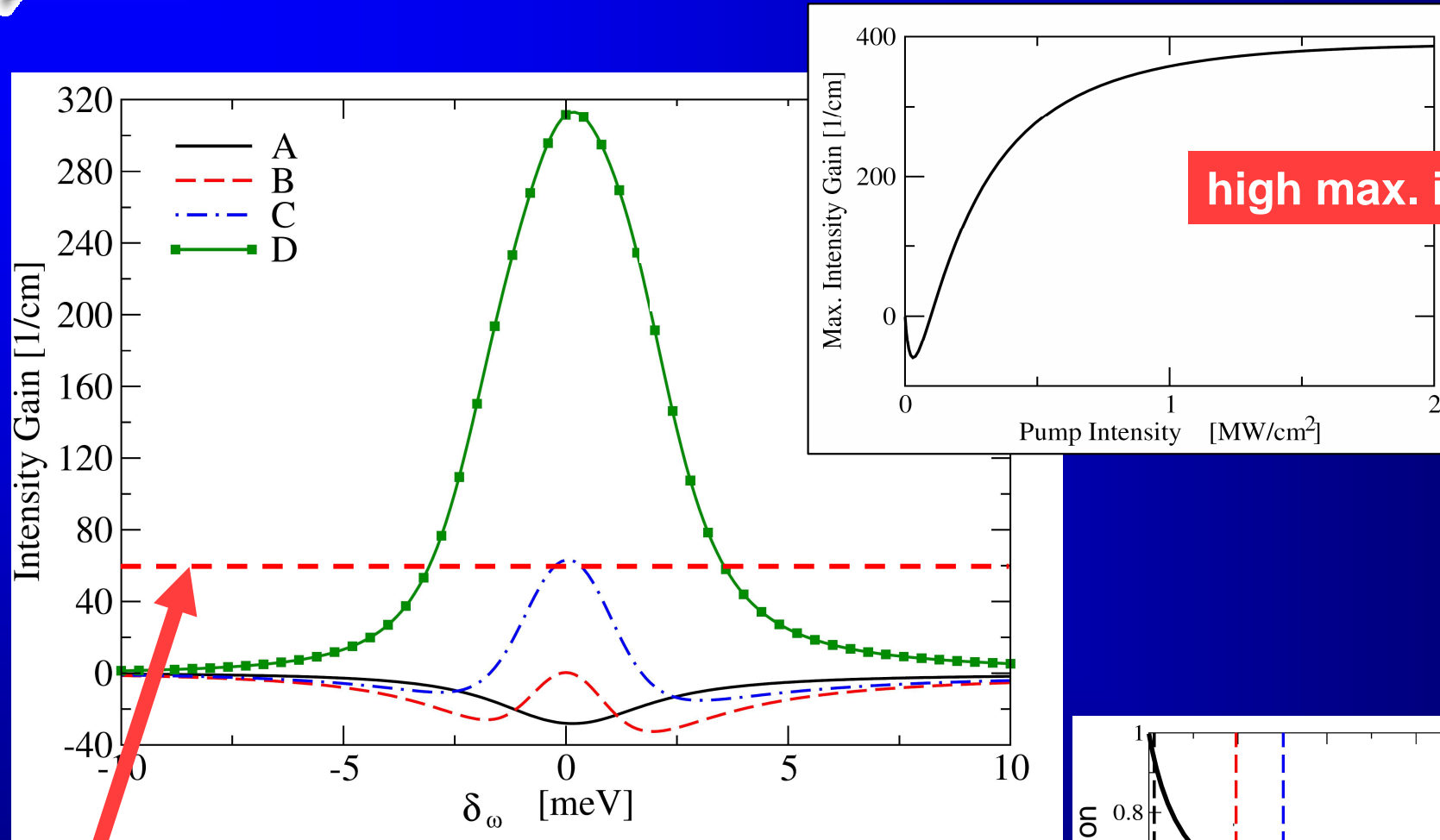


$$\alpha_{THz} = \alpha_{WG} + \alpha_M + \alpha_{FCL} \geq 60 \text{ cm}^{-1}$$

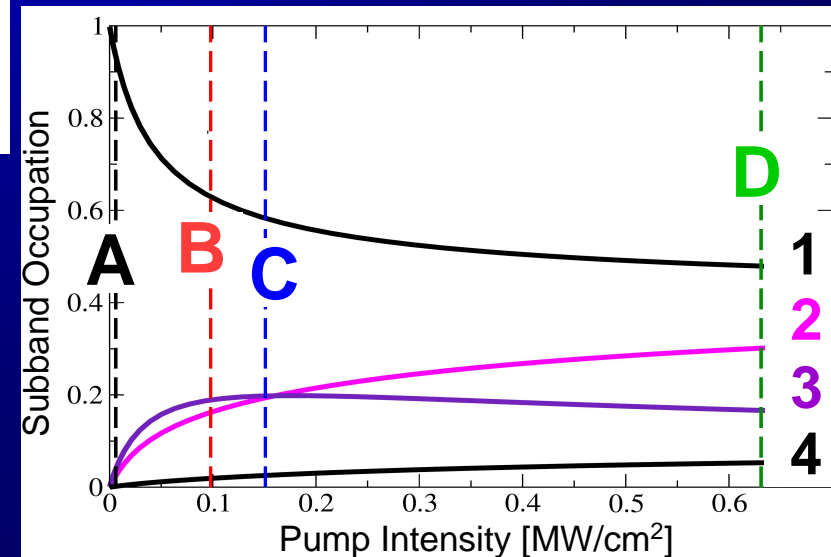
SIGN OF QUANTUM COHERENCE!!!



Let's look at the linear gain for different drive intensities

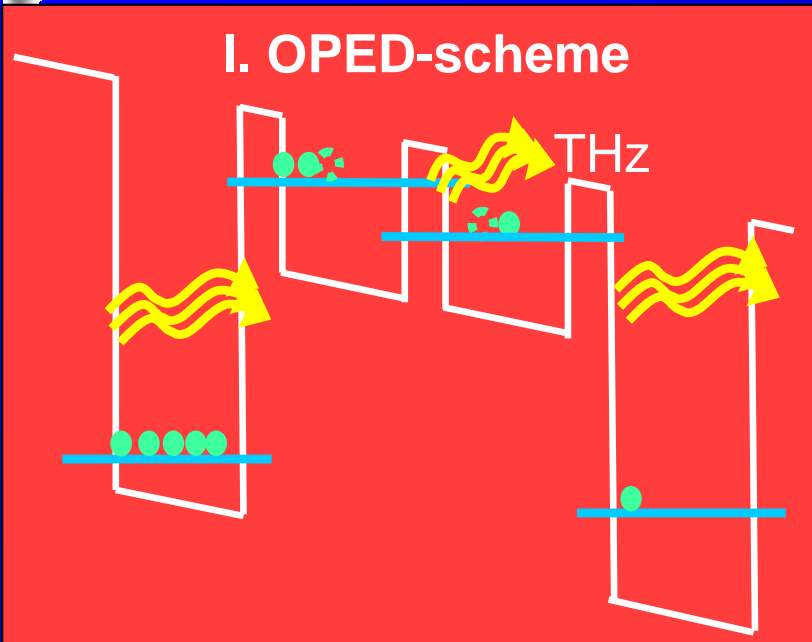


$$\alpha_{THz} = \alpha_{WG} + \alpha_M + \alpha_{FCL} \geq 60 \text{ cm}^{-1}$$

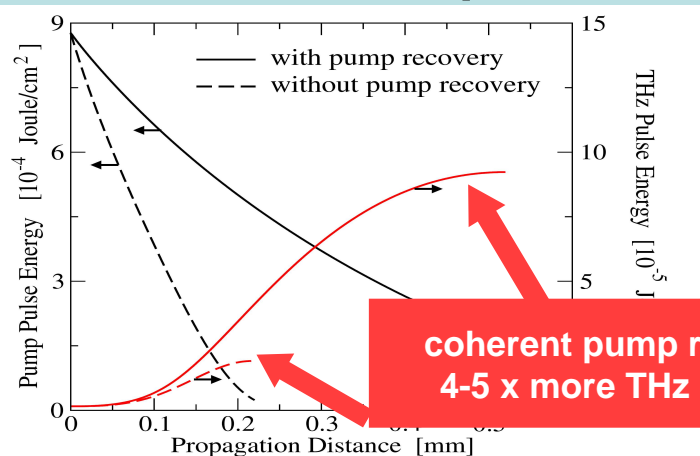


Summary: more information in PRL (Sept. 2007)

I. OPED-scheme

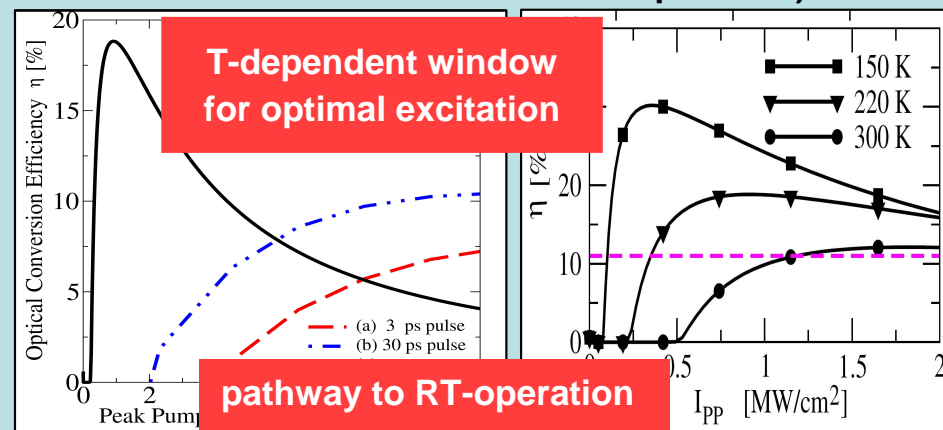


II. Benefits of Pump Recovery



III. Capabilities of OPED

(Dependence of Optical Conversion Efficiency on Pulse Duration and Temperature)



IV. Dependence of Small Signal Gain on Pump Intensity

